

Dynamics in dissipative quantum systems

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Motivation

"The simplest model involving quantum coherence is a two-state system.[...] In reality, such a system is strongly affected by the surroundings"

– Quantum dissipative systems, U. Weiss

how we use "coherence"

-in a small quantum system

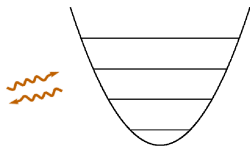
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how we use "coherence"

- in a small quantum system coupled to a large environment



- spin-boson model: widely used in physics, chemistry and quantum information sciences

- here: unbiased, ohmic spin boson model in scaling limit

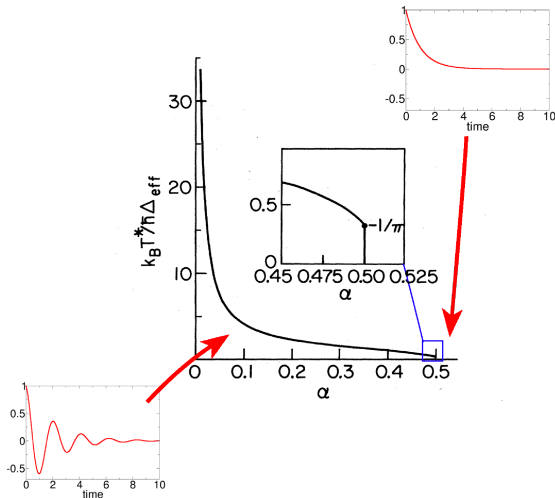
Goal and Approach

- dissipative system: relaxation dynamics in coupling(α)- T -plane

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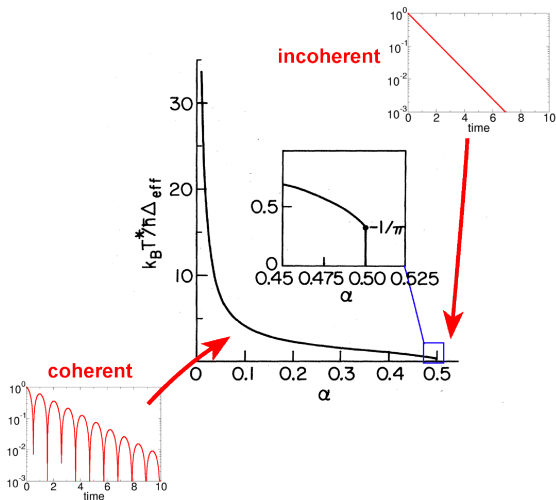
Garg '85; Weiss, Grabert '86; Leggett et al. '87



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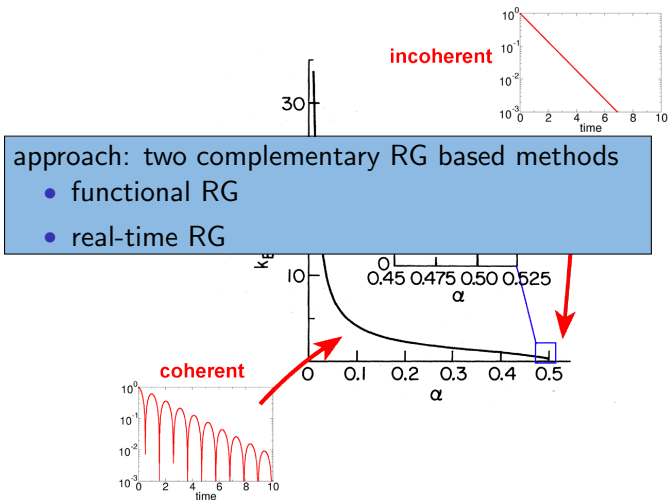
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Model

Hamiltonian

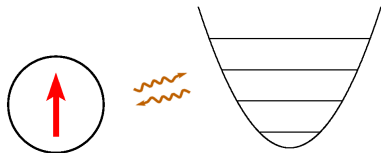
$$H = -\frac{\Delta}{2}\sigma_x - \sum_k \frac{\lambda_k}{2}\sigma_z (b_k^\dagger + b_k) + \sum_k \omega_k b_k^\dagger b_k$$

coupling: spectral density

$$J(\omega) = \sum_k \lambda_k^2 \delta(\omega - \omega_k)$$

ohmic case: rich physics

$$J(\omega) = 2\alpha\omega\Theta(\omega_c - \omega)$$



initial state: spin-up \times therm. eq.

Model

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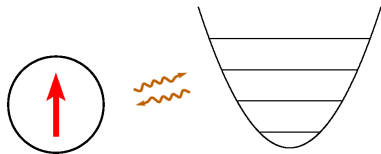
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scaling limit: $\omega_c \rightarrow \infty$

emergent (Kondo) scale: $T_K = \Delta(\Delta/\omega_c)^{\alpha/(1-\alpha)}$

Model: $|\alpha - 1/2| \ll 1$

dynamics: coherent (osz. damped) \leftrightarrow incoherent (monotonic)

spin expect. value $P(t) = \langle \sigma_z(t) \rangle$: $\alpha < 1/2$ coh., $1/2 < \alpha < 1$ incoh.

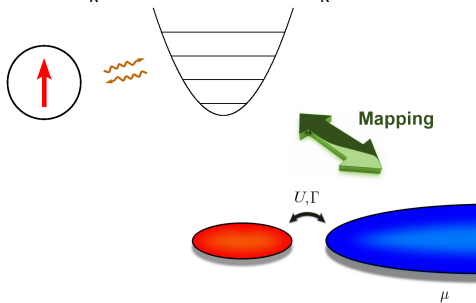
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bosonization: maps to interacting resonant level model (IRLM)

$$H = \sum_k \epsilon_k a_k^\dagger a_k + \sqrt{\frac{\Gamma_0}{2\pi\nu}} \sum_k (d^\dagger a_k + d a_k^\dagger) + \frac{U}{2\nu} (d^\dagger d - d d^\dagger) \sum_{kk'} : a_k^\dagger a_{k'} :$$



\rightarrow mapping of parameters:

$$\Gamma_0 = \Delta^2 / \omega_c$$

$$U = 1 - \sqrt{2\alpha}$$

$$g = 2U - U^2 = 1 - 2\alpha$$

$$P(t) = \langle \sigma_z(t) \rangle = 1 - 2n(t)$$

$n(t)$: occupancy of level

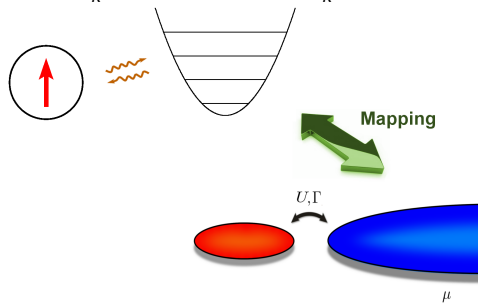
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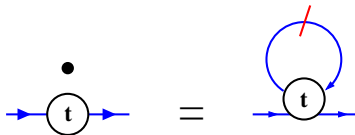
- **ADVANTAGE:** $\alpha = 1/2 \rightarrow g = U = 0$ (Toulouse limit)

Methods: FRG and RTRG

Functional RG: Keldysh Space

two-particle interaction: FRG

- cut-off Λ in Green functions:
 $G^0 \rightarrow G^{0,\Lambda}$ with $G^{0,\Lambda \rightarrow \infty} \rightarrow 0$ and $G^{0,\Lambda=0} \rightarrow G^0$
- hierarchy of flow equations for vertex functions
- truncate hierarchy after first order (only Σ flows)



$$\partial_\Lambda \gamma_1^\Lambda(\mathbb{1}, \mathbb{1}') = \sum_{\mathbb{2}\mathbb{2}'} \hat{S}_{\mathbb{2},\mathbb{2}'}^\Lambda \gamma_2^\Lambda(\mathbb{1}\mathbb{2}; \mathbb{1}'\mathbb{2}') \quad (\text{with: } \mathbb{1} = \{t, i, p\})$$

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$$\hat{S}_{11}^\Lambda = S^\Lambda(\mathbb{I}, \mathbb{I}') = \sum_{22'} G_{1,2'}^\Lambda \left[\partial_\Lambda [G^{0,\Lambda}]^{-1} \right]_{2'2} G_{21'}^\Lambda = -\partial_\Lambda^* \hat{G}_{11}^\Lambda$$

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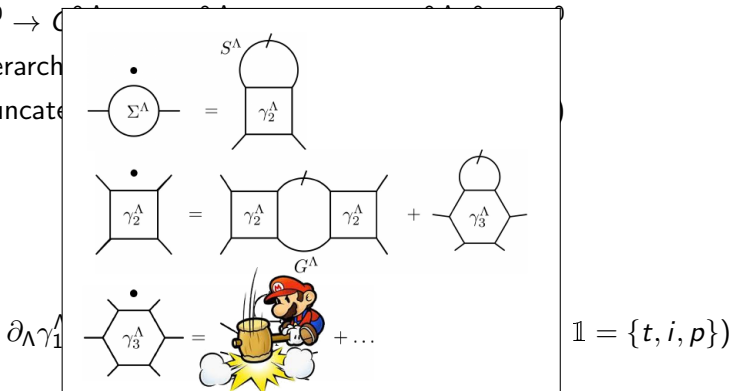
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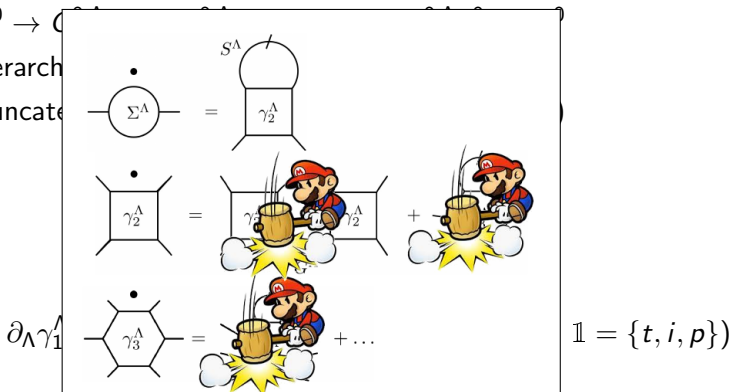
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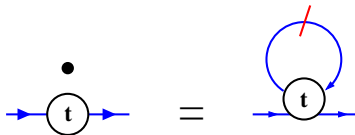


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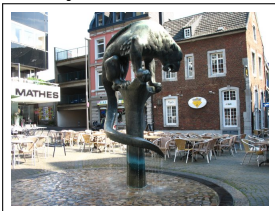


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Functional RG: Flexibility

steady driven state



periodic driving

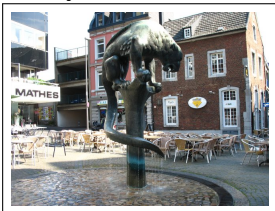


transient dynamics



Functional RG: Flexibility

steady driven state



$t_0 \rightarrow -\infty$:

$$G(t, t') = G(t - t')$$

Fouriertransform
 $\Rightarrow G(\omega)$

periodic driving

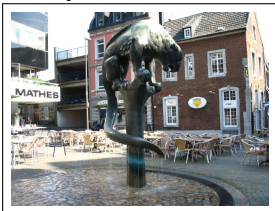


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$t_0 \rightarrow -\infty$:

$$G(t, t') = G(t + T, t' + T)$$

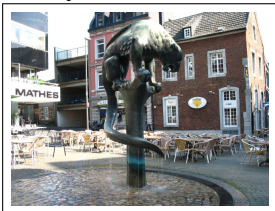
Floquet theory $\Rightarrow G_k(\omega)$

transient dynamics



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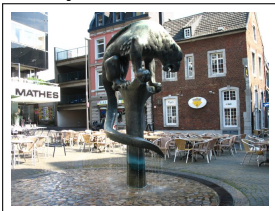


finite t_0 :

keep $G(t, t')$ explicit

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Real-Time RG: Liouville-Laplace Space

aim at reduced density matrix of local system

$$i\dot{\rho}(t) = [H, \rho] = L\rho$$

$$\rho(t) = e^{-iL(t-t_0)}\rho(t_0)$$

$$\rho_S(t) = \text{Tr}_{\text{res}}[\rho(t)]$$

$$\rho_S(t) = P(t)\rho_S(t_0) = \frac{i}{2\pi} \int_{-\infty+i0^+}^{\infty+i0^+} dE e^{-iEt} \Pi_1(E)\rho_S(t_0),$$

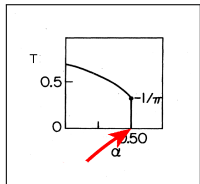
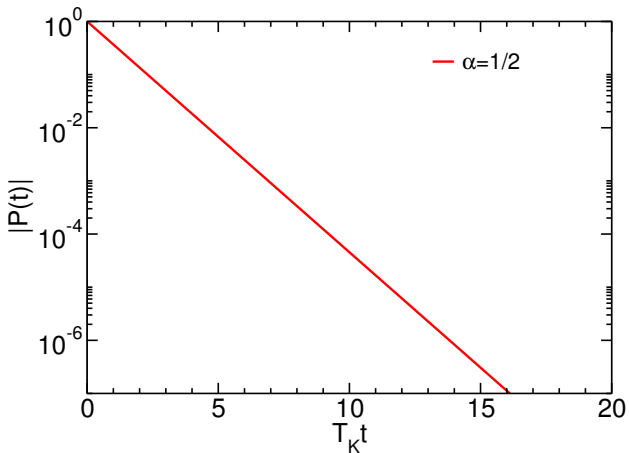
with $\Pi_1(E) = [E + i\Gamma_1(E)]^{-1}$ and Laplace variable E as flow parameter:

$$\frac{d\Gamma_{1/2}(E)}{dE} = -(1 - 2\alpha)\Gamma_1(E)\Pi_{2/1}(E),$$

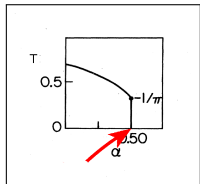
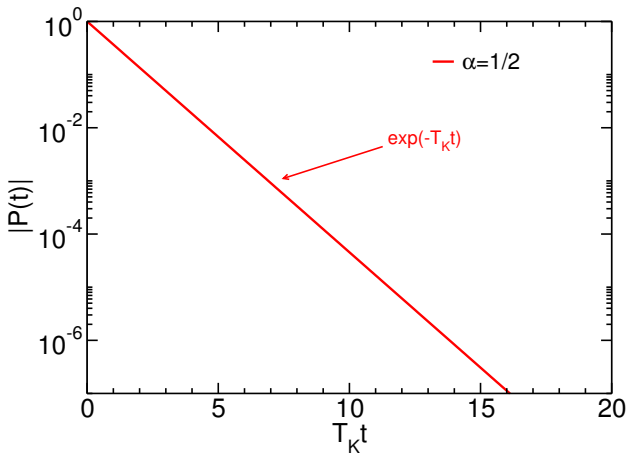
with $\Pi_2(E) = [E + i\Gamma_2(E)/2]^{-1}$ and initially $\Gamma_n(i\omega_c) = \Delta^2/\omega_c$

Dynamics in the α -T-plane

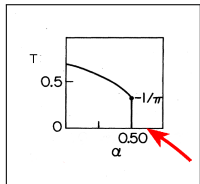
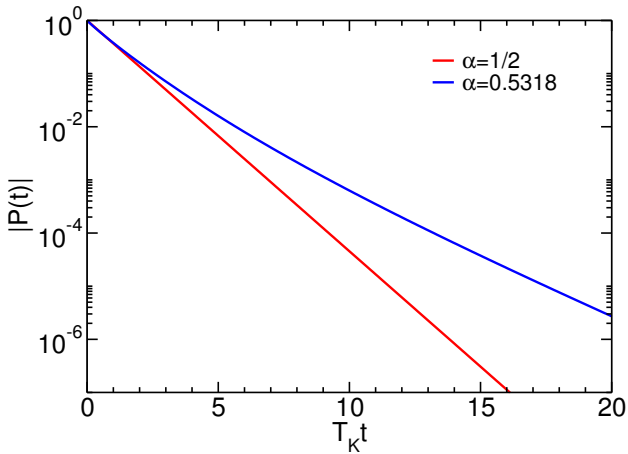
Spin Relaxation Dynamics: $T = 0$



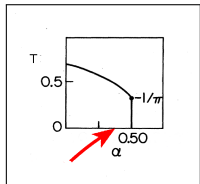
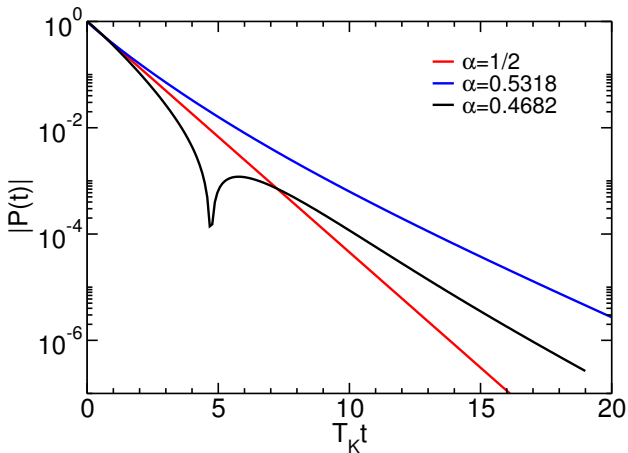
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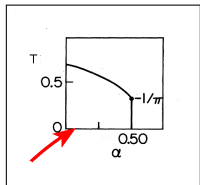
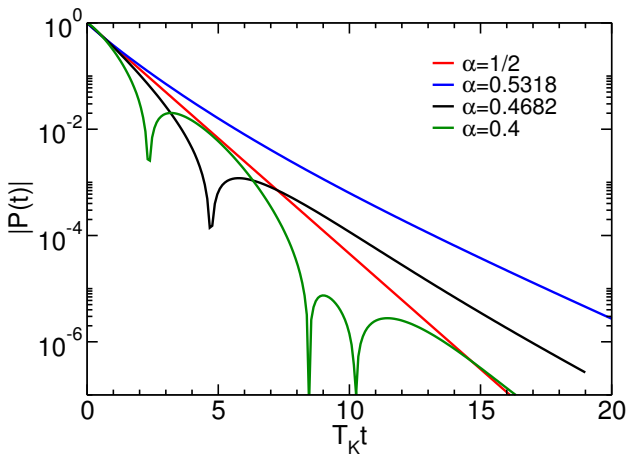
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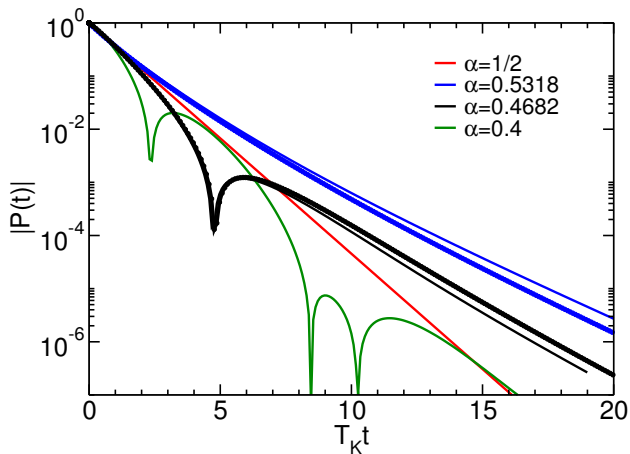
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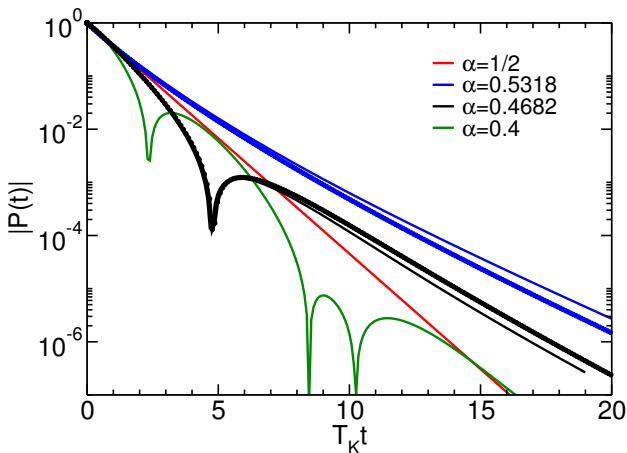


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RTRG!

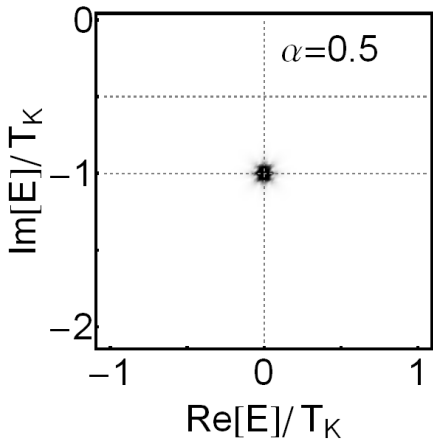
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coherent regime: non-monotonic only for small-intermediate $T_K t$?

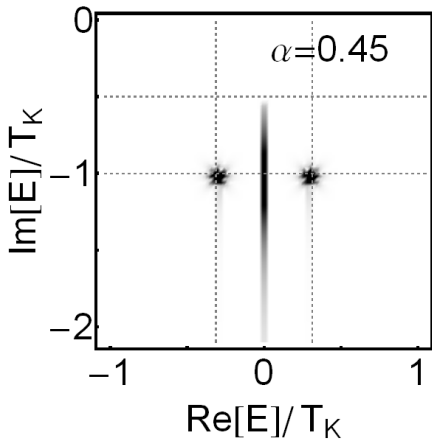
RTRG: Numerical Solution in Complex E -Plane

$$|\partial_{\text{Im}E} \text{Re} [\Pi_1] + \partial_{\text{Re}E} \text{Im} [\Pi_1]|$$



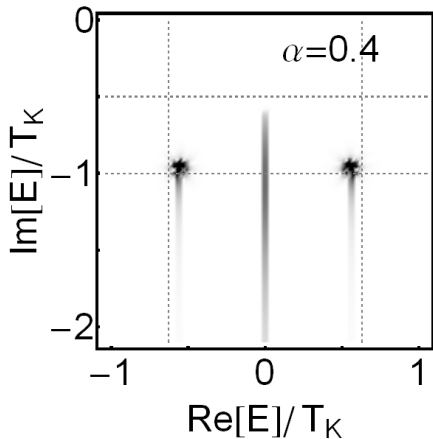
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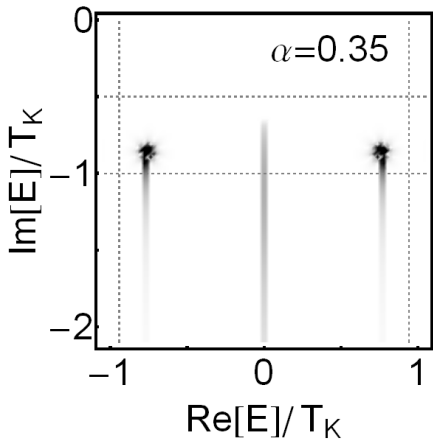
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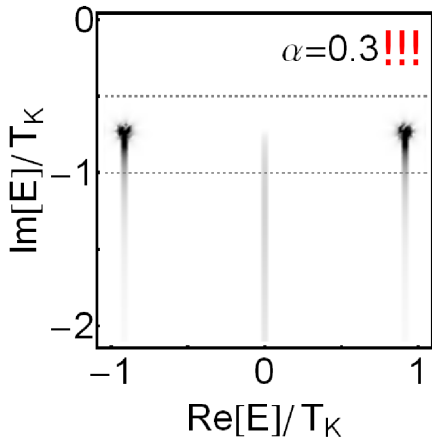
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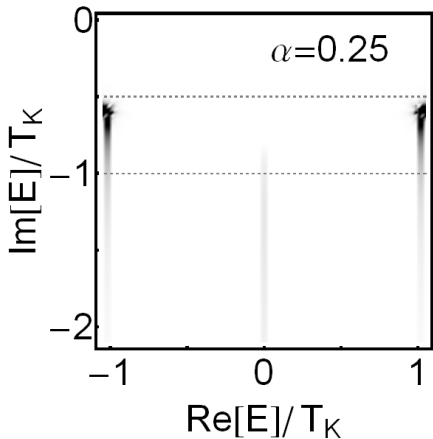
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Intermezzo: Conclusion

- three distinct dynamical regimes
 - $1/2 < \alpha < 1$: incoherent
 - $\alpha_c < \alpha < 1/2$: partially coherent
 - $0 < \alpha < \alpha_c$: asymptotically coherent

with $\alpha_c \approx 0.3$

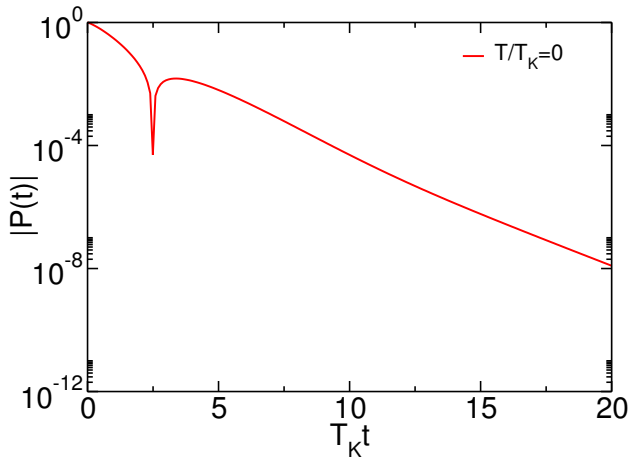
- still controlled?

RTRG for SBM (small α !) gives $\alpha_c \approx 0.36$

(Kashuba, Schoeller '13)

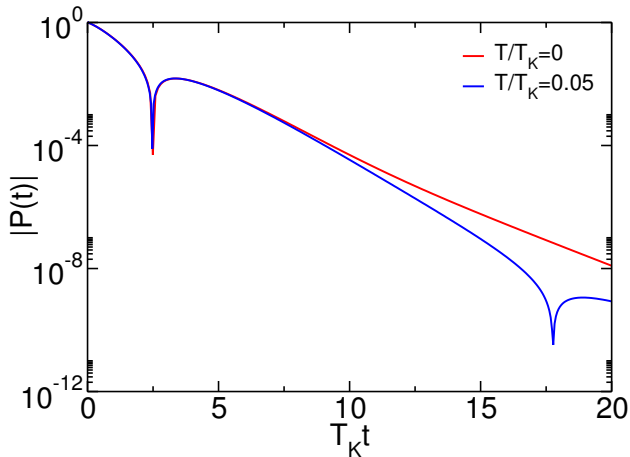
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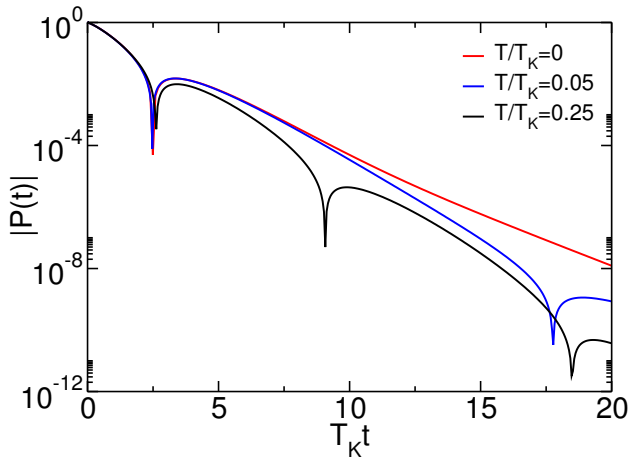
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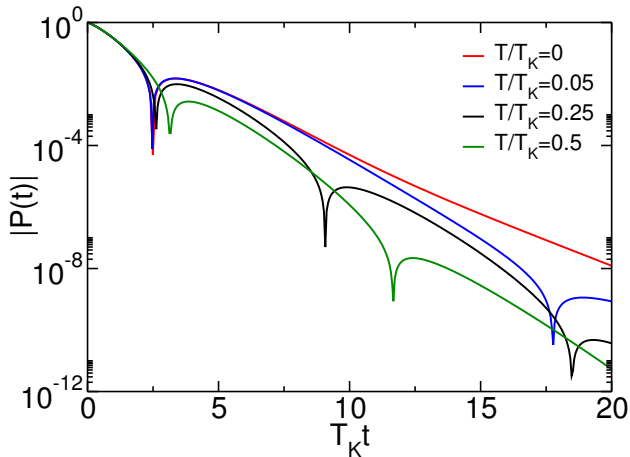
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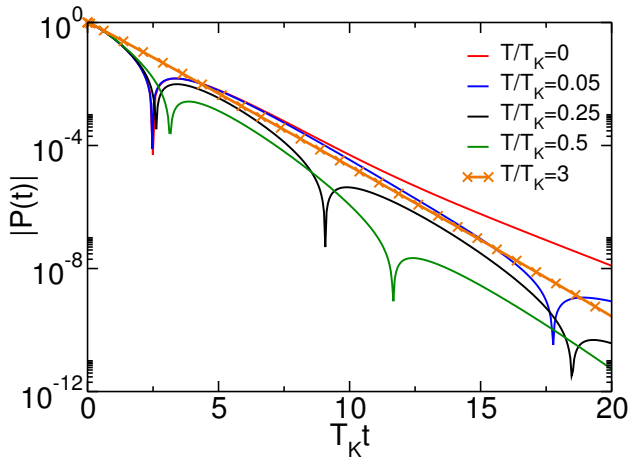
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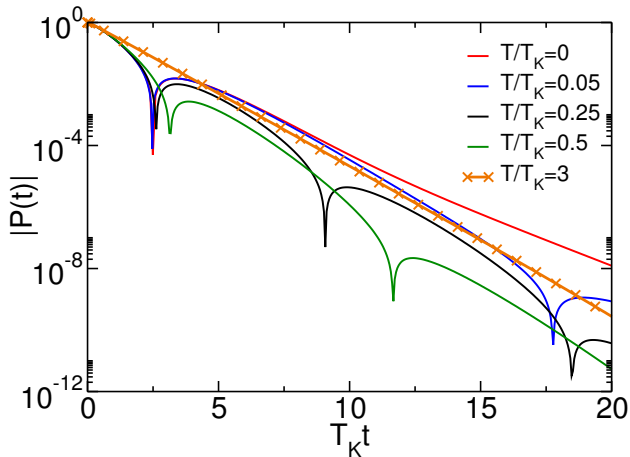
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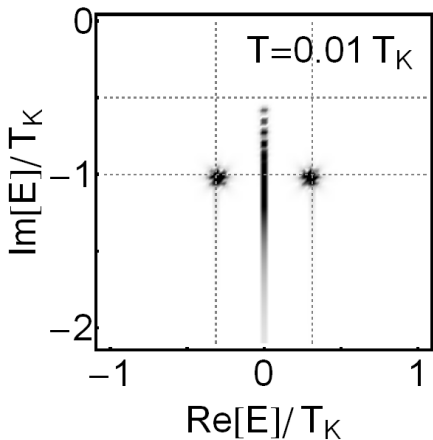
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- partially coherent \rightarrow asymptotically coherent \rightarrow incoherent?

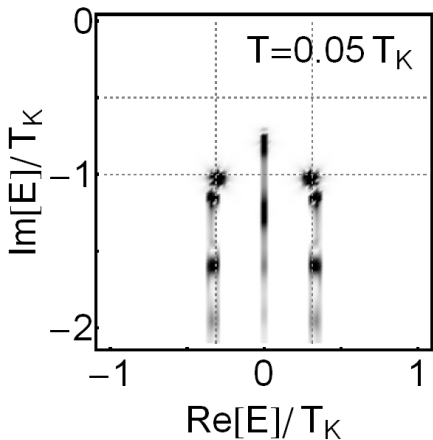
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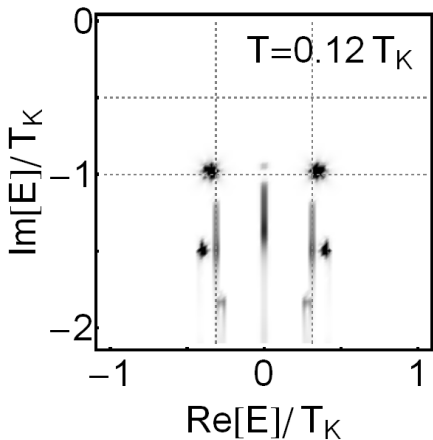
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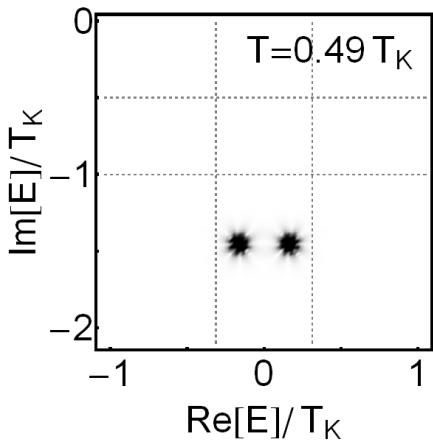
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$T_{c_1}(\alpha)!!!$

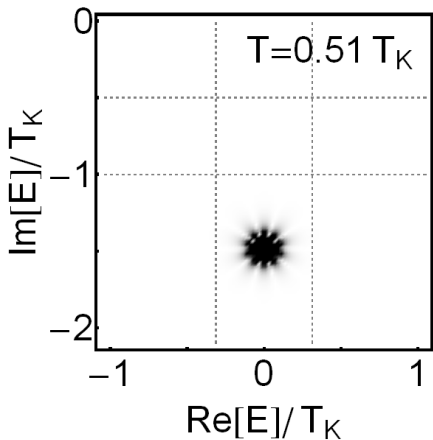
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RTRG: Numerical Solution in Complex E -Plane $\alpha = 0.45$

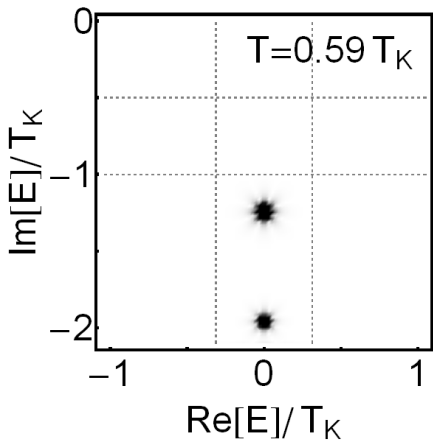
$$|\partial_{\text{Im}E} \text{Re} [\Pi_1] + \partial_{\text{Re}E} \text{Im} [\Pi_1]|$$



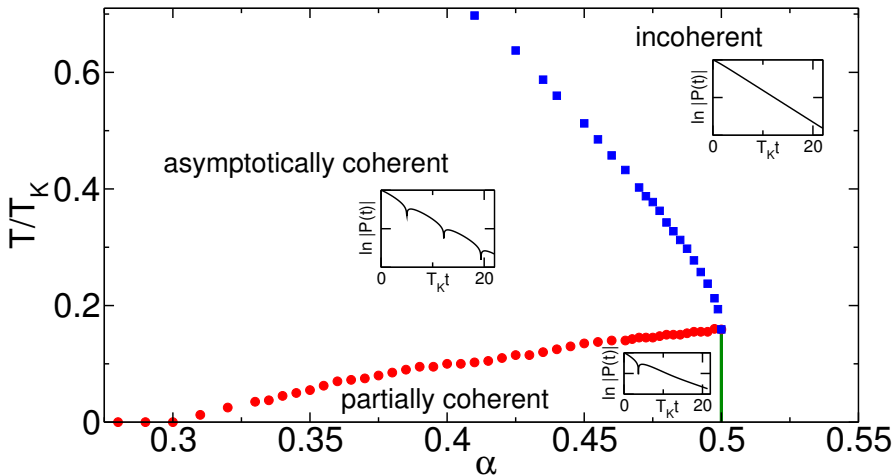
$T_c(\alpha)!!!$

RTRG: Numerical Solution in Complex E -Plane $\alpha = 0.45$

$$|\partial_{\text{Im}E} \text{Re} [\Pi_1] + \partial_{\text{Re}E} \text{Im} [\Pi_1]|$$



Conclusion: Dynamical Regimes



$T = 0$: PRL **110**, 100405 (2013)
 $T > 0$: PRB **88**, 241110(R) (2013)

Summary

dynamics in α - T -plane

- three distinct dynamical regimes
- increase T : transition from partially to asymptotically coherent
- 'coherence from elevated temperature'

$T = 0$: Phys. Rev. Lett. **110**, 100405 (2013)

Quench: Phys. Rev. B **88**, 165133 (2013)

$T > 0$: Phys. Rev. B **88**, 241110(R) (2013)

avenues of future research

- periodic driving (see Eissing et al. arXiv:1508.01325)
- current noise
- ...