

Trogir (Croatia), 18 - Sept - 2015

Superconductivity in nanoscopic systems

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<http://kft.umcs.lublin.pl/doman/lectures>

Outline

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A few questions:

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⇒ **how can we obtain nano-superconductivity**

/ proximity effect /

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/ spectroscopic signatures /

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⇒ **how can we obtain nano-superconductivity**

/ proximity effect /

⇒ **how can we observe nano-superconductivity**

/ spectroscopic signatures /

⇒ **where can we use nano-superconductors**

/ practical aspects /

1. Nano-superconductivity:

⇒ **how to obtain it ?**

Superconducting state

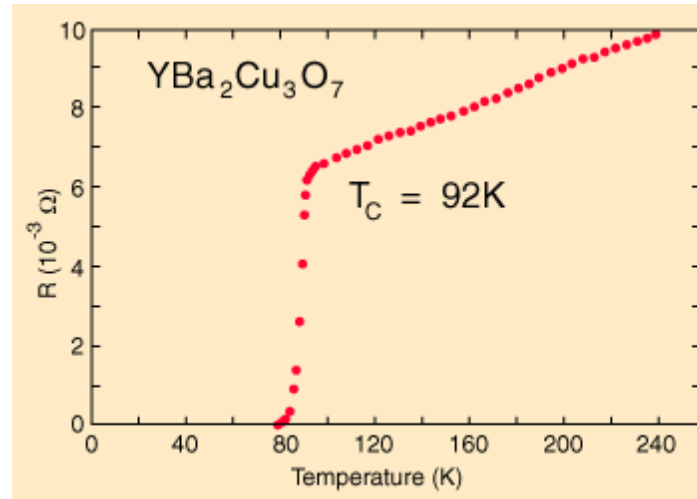
– of bulk materials

Superconducting state

– of bulk materials



ideal d.c. conductance

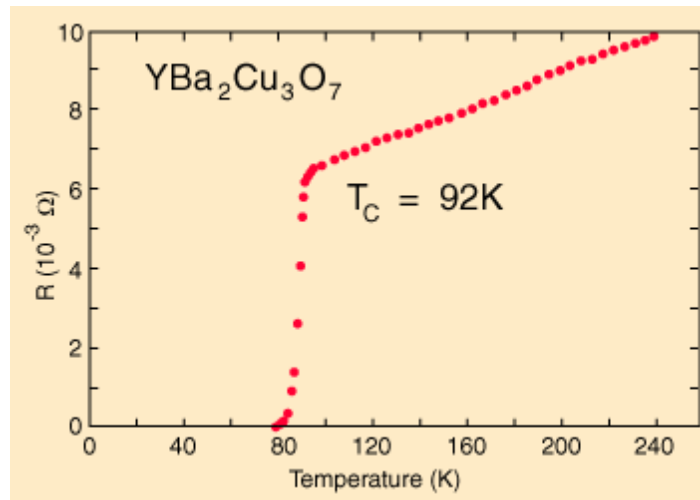


Superconducting state

– of bulk materials

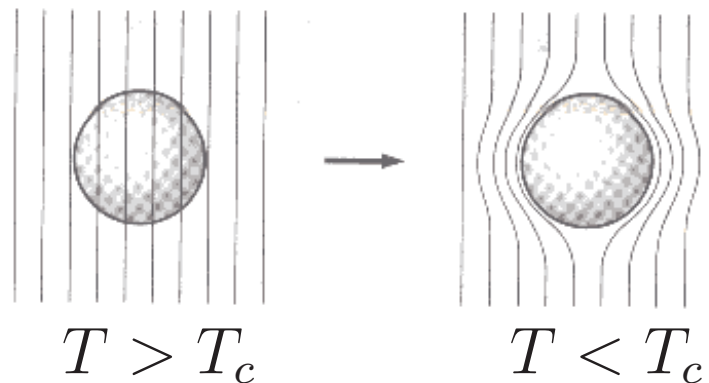


ideal d.c. conductance



ideal diamagnetism

/perfect screening of the d.c. magnetic field/



Superconducting state – of bulk materials

⇒ ideal d.c. conductance (vanishing resistance)

⇒ ideal diamagnetism (Meissner effect)

are caused by the superfluid electron pairs

$$n_s(T) \propto 1/\lambda^2$$

that move coherently over macroscopic distances.

Superconducting state – of bulk materials

The pairing mechanisms originate from:

1. **phonon-exchange**

/ classical superconductors, MgB₂, ... /

2. **magnon-exchange**

/ heavy fermion compounds /

3. **strong correlations**

/ spin exchange $\frac{2t_{ij}^2}{U}$ in the high T_c superconductors /

.. **other exotic processes**

/ ultracold atoms, nuclei, gluon-quark plasma /

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Onset of the fermion pairing often goes hand in hand with appearance of the **superconductivity/superfluidity**, but it doesn't have to be a rule.

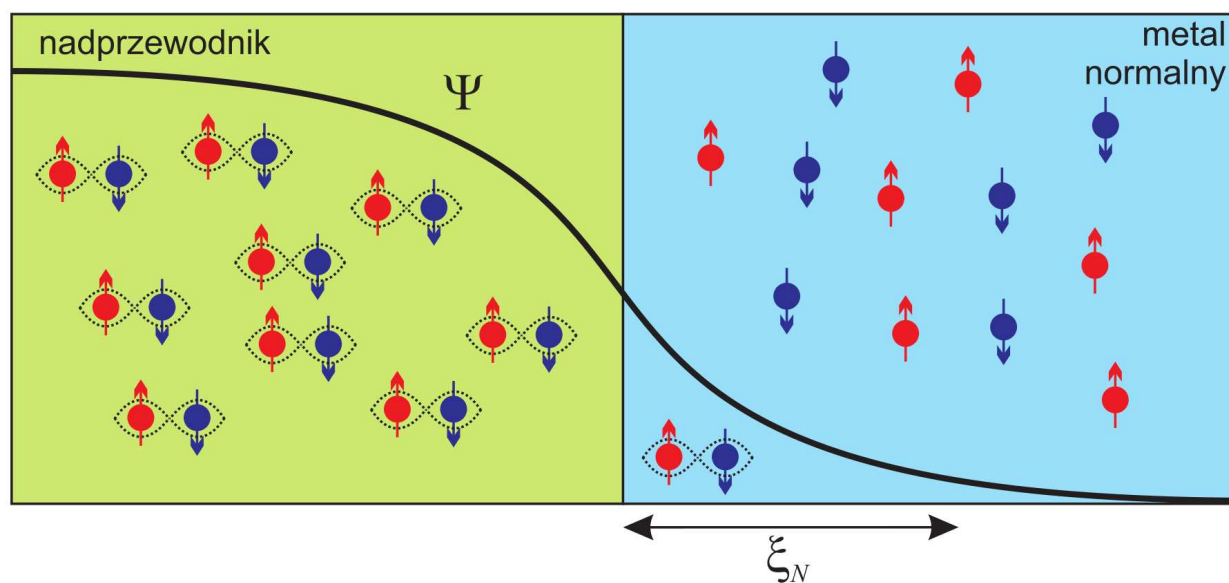
Proximity effect

– induced superconductivity

Proximity effect

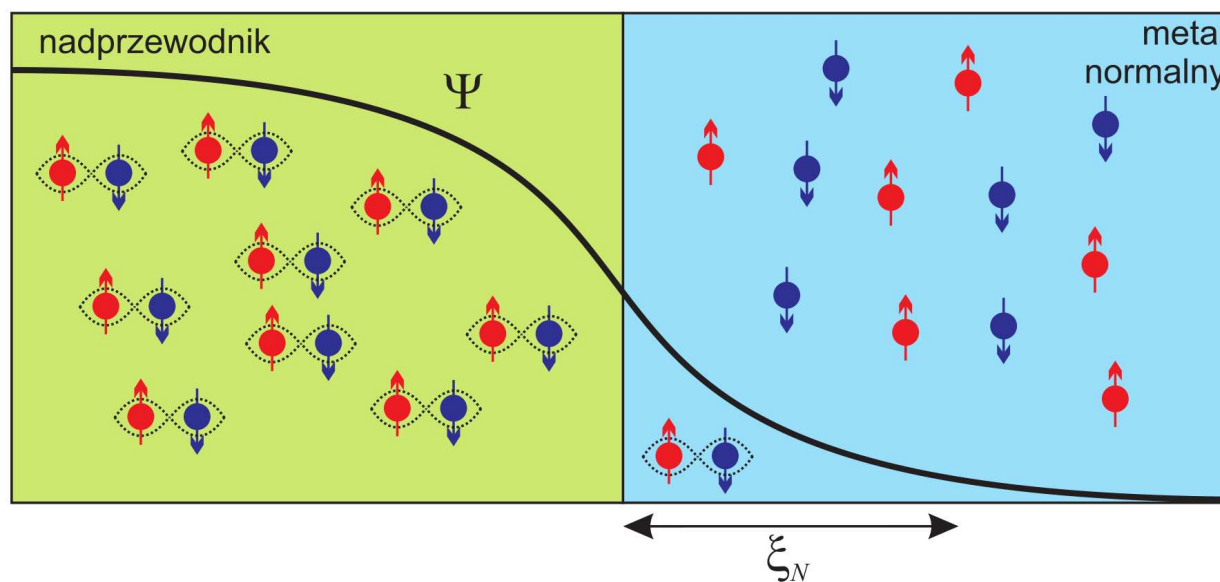
– induced superconductivity

★ Any material brought in contact with superconductor



Proximity effect – induced superconductivity

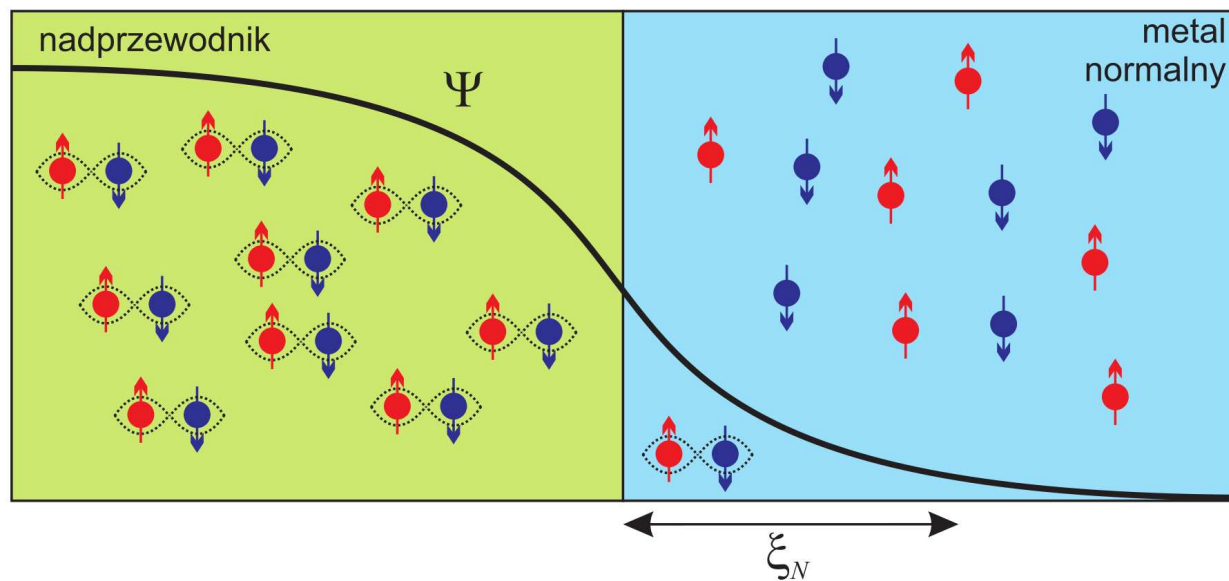
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absorbs the paired electrons up to distances $\sim \xi_n$.

Proximity effect – induced superconductivity

★ Any material brought in contact with superconductor



absorbs the paired electrons up to distances $\sim \xi_n$.

★ Spatial size L of nanoscopic objects is $L \ll \xi_n$!

2. Nano-superconductivity:

⇒ how can we observe it ?

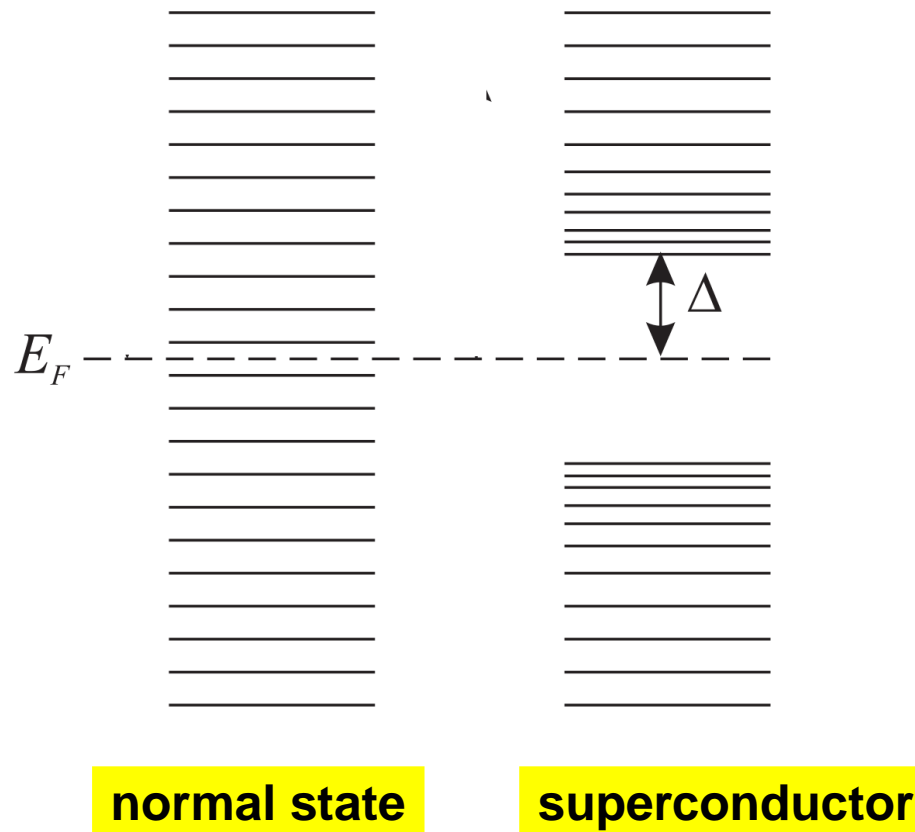
Superconductivity in nanosystems

– **specific issues**

Superconductivity in nanosystems

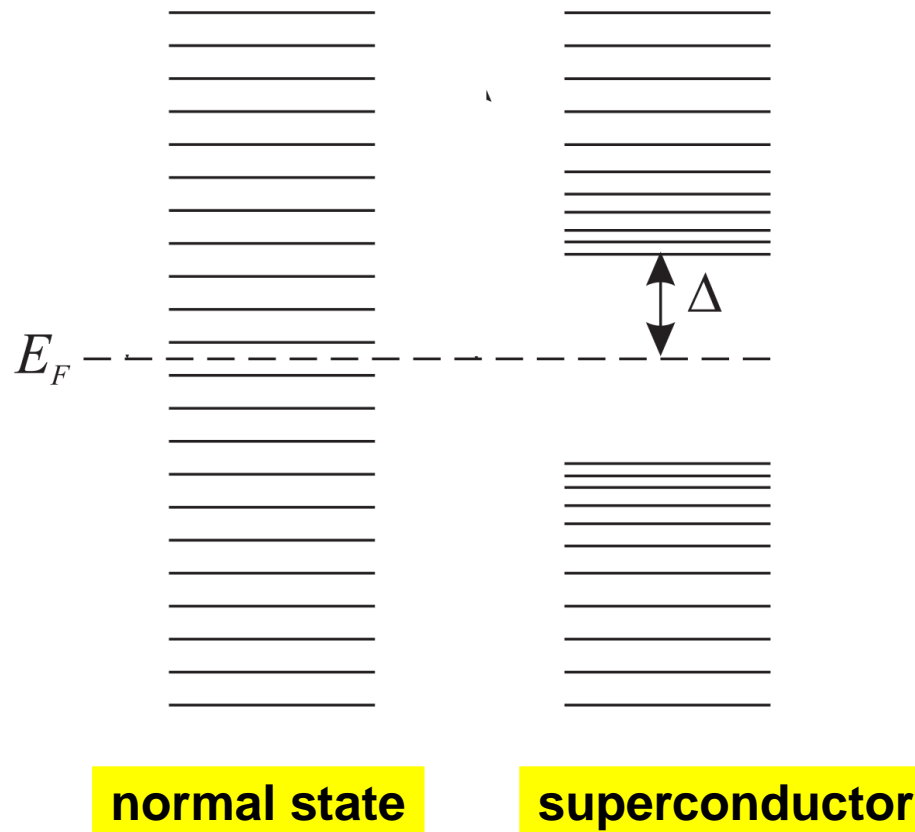
– specific issues

1. Quantum Size Effect \longrightarrow discrete energy spectrum



Superconductivity in nanosystems – specific issues

1. Quantum Size Effect \longrightarrow discrete energy spectrum



Anderson criterion:

superconductivity only for $\Delta > \varepsilon_{i+1} - \varepsilon_i$

Superconductivity in nanosystems – **specific issues**

2. Coulomb blockade (electron pairing vs repulsion)

The odd / even electron number plays the important and qualitative role !

Superconductivity in nanosystems – specific issues

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Coulomb potential U_C is usually much smaller than Δ , therefore its influence can be in practice observed only indirectly, via :

$$|\uparrow\rangle \iff u|0\rangle - v|\uparrow\downarrow\rangle$$

(quantum phase transition)

Superconductivity in nanosystems – specific issues

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Physical consequences:

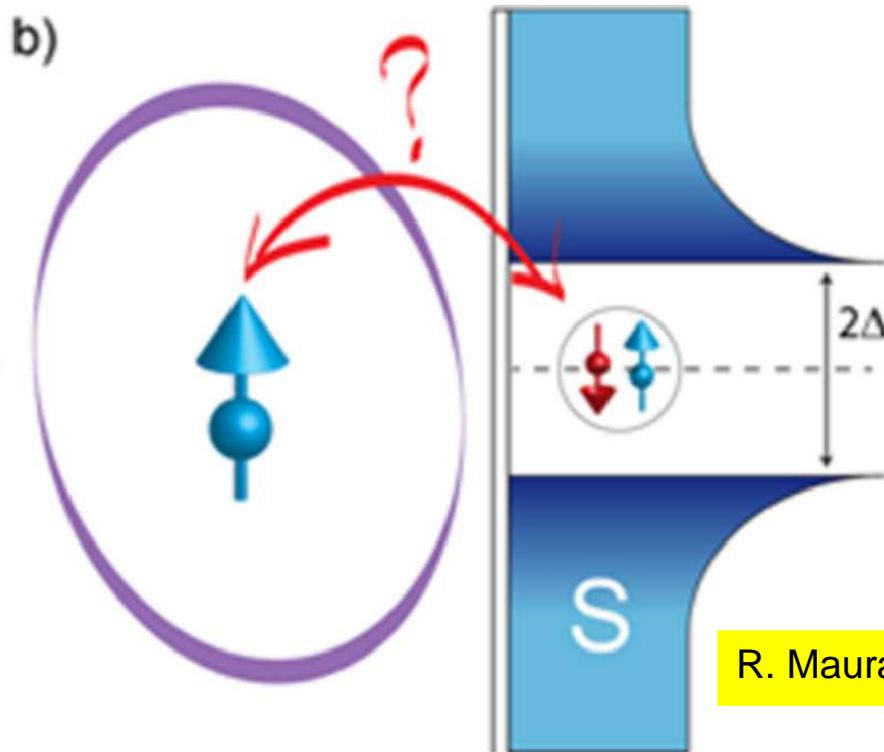
⇒ inversion of the Josephson current (in S-QD-S junctions)

⇒ activation/blocking of the Kondo effect (in N-QD-S junctions)

Superconductivity in nanosystems

– specific issues

3. Pairing vs Kondo state ('to screen or not to screen')

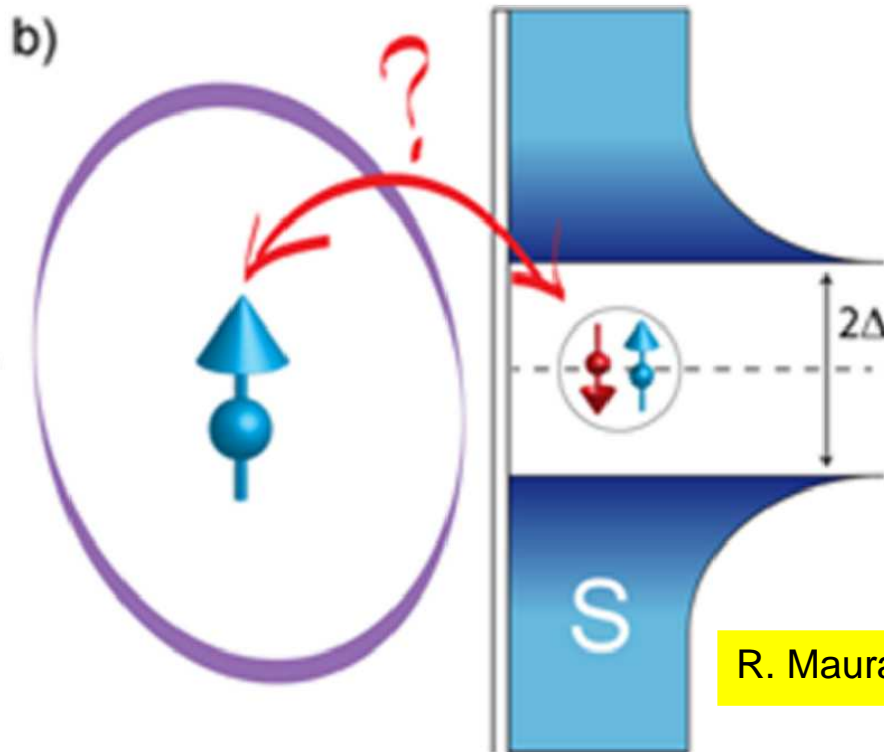


R. Maurand, Ch. Schönberger, *Physics* **6**, 75 (2013).

Superconductivity in nanosystems

– specific issues

3. Pairing vs Kondo state ('to screen or not to screen')



R. Maurand, Ch. Schönberger, Physics **6**, 75 (2013).

⇒ **states near the Fermi level are depleted**

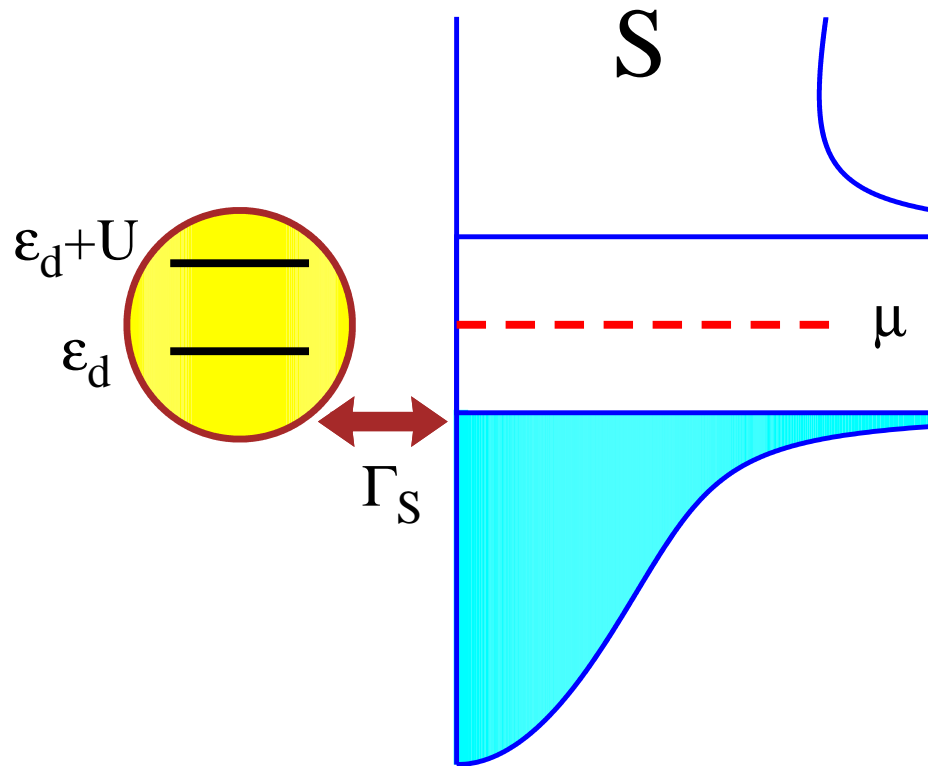
⇒ **electron pairing vs the Kondo state (nontrivial relation)**

Theoretical model

– **single Anderson impurity**

Theoretical model – single Anderson impurity

The single quantum impurity (dot) coupled to superconducting reservoir



ϵ_d – energy level, U – Coulomb potential, Γ_S – hybridization

Theoretical model – single Anderson impurity

Hamiltonian

$$\begin{aligned}\hat{H} &= \sum_{\sigma} \epsilon_d \hat{d}_{\sigma}^{\dagger} \hat{d}_{\sigma} + U \hat{n}_{d\uparrow} \hat{n}_{d\downarrow} \\ &+ \sum_{\mathbf{k}, \sigma} \left(V_{\mathbf{k}} \hat{d}_{\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} + V_{\mathbf{k}}^* \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{d}_{\sigma} \right) + \hat{H}_S\end{aligned}$$

where

$$\hat{H}_S = \sum_{\mathbf{k}, \sigma} (\epsilon_{\mathbf{k}} - \mu) \hat{c}_{\mathbf{k}\sigma}^{\dagger} \hat{c}_{\mathbf{k}\sigma} - \sum_{\mathbf{k}} \left(\Delta \hat{c}_{\mathbf{k}\uparrow}^{\dagger} \hat{c}_{\mathbf{k}\downarrow}^{\dagger} + \text{h.c.} \right)$$

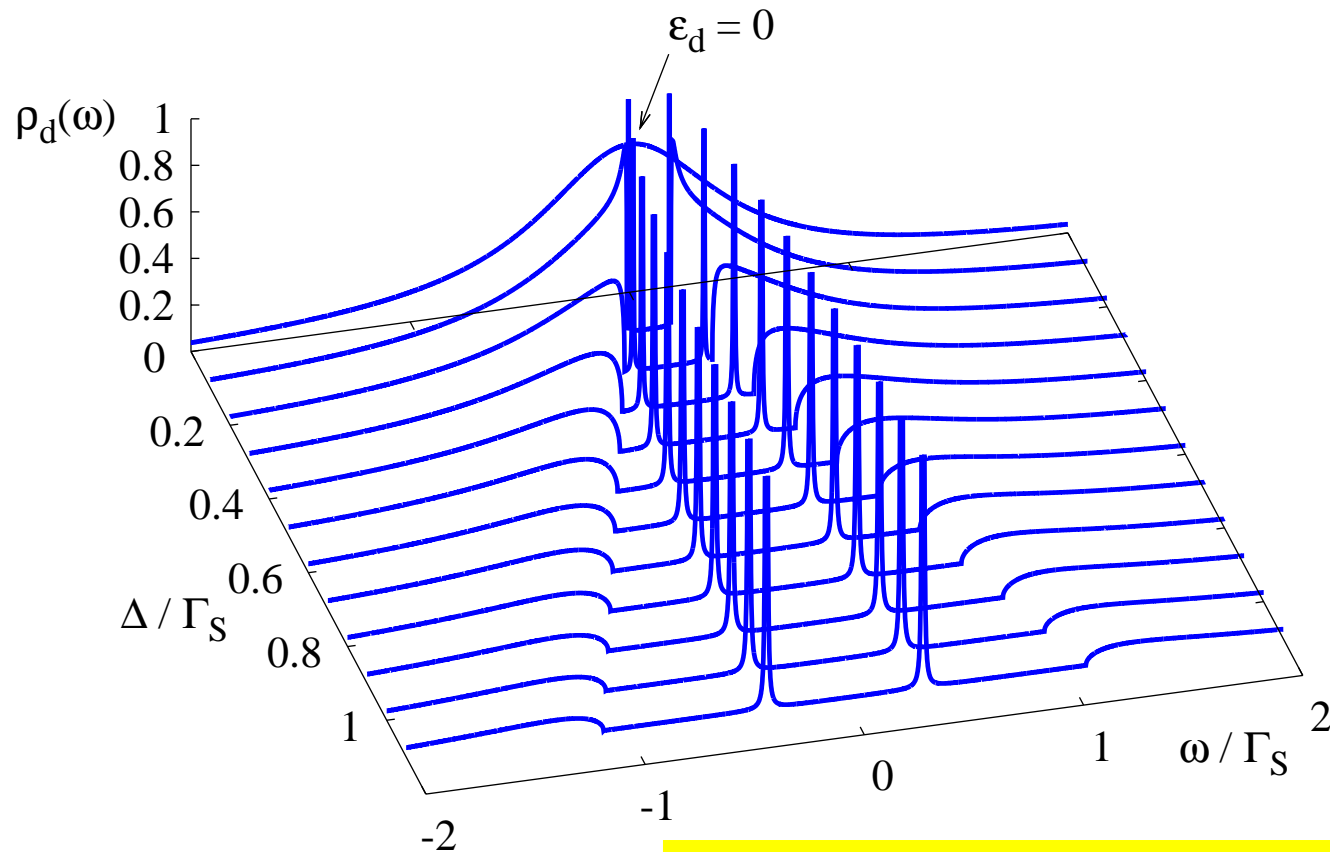
describes the BCS-type superconductor.

Microscopic description

- exact solution for $U = 0$

Microscopic description

– exact solution for $U = 0$



J. Phys.: Condens. Matter **25**, 435305 (2013).

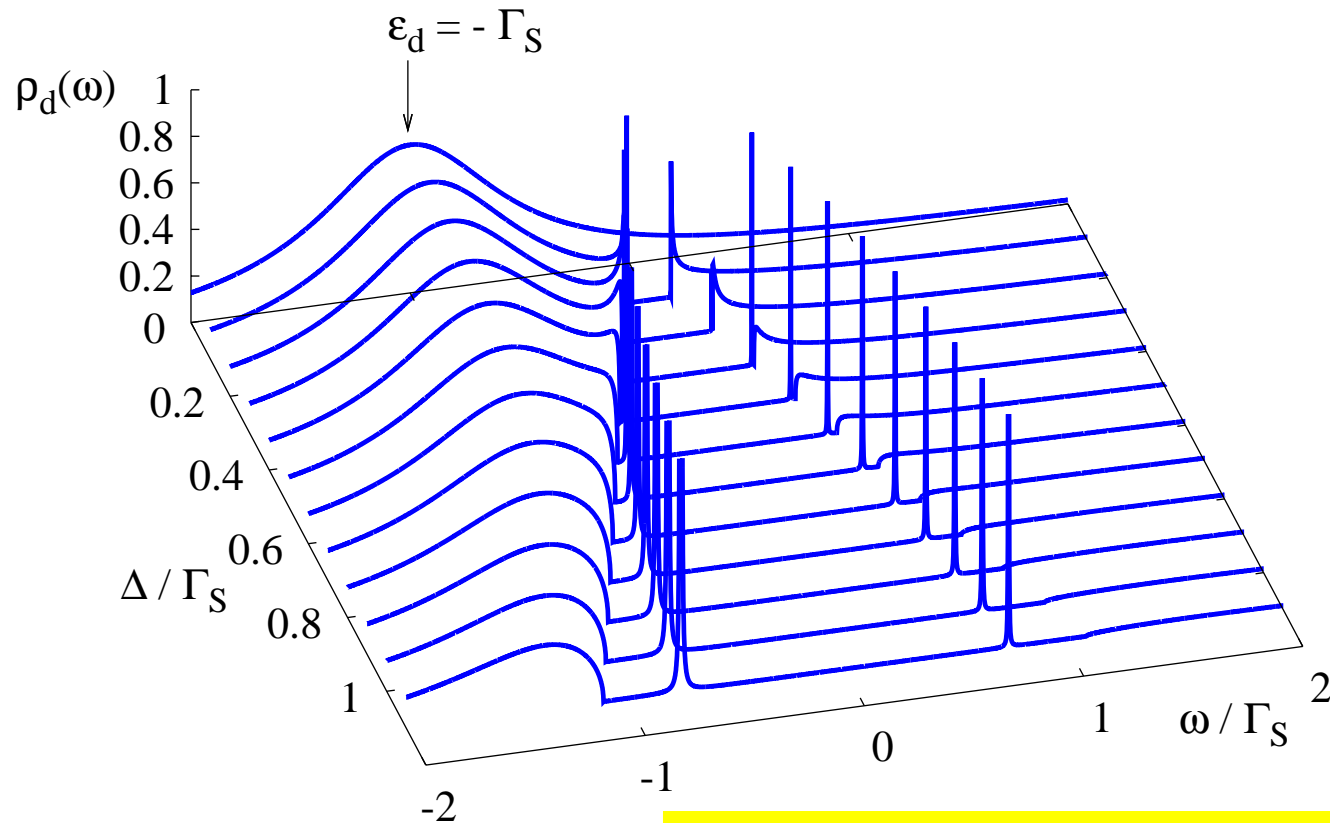
Spectrum consists of:

⇒ a continuum at energies $-\Delta < \omega < \Delta$

⇒ in-gap resonances (Andreev bound states)

Microscopic description

– exact solution for $U = 0$



J. Phys.: Condens. Matter **25**, 435305 (2013).

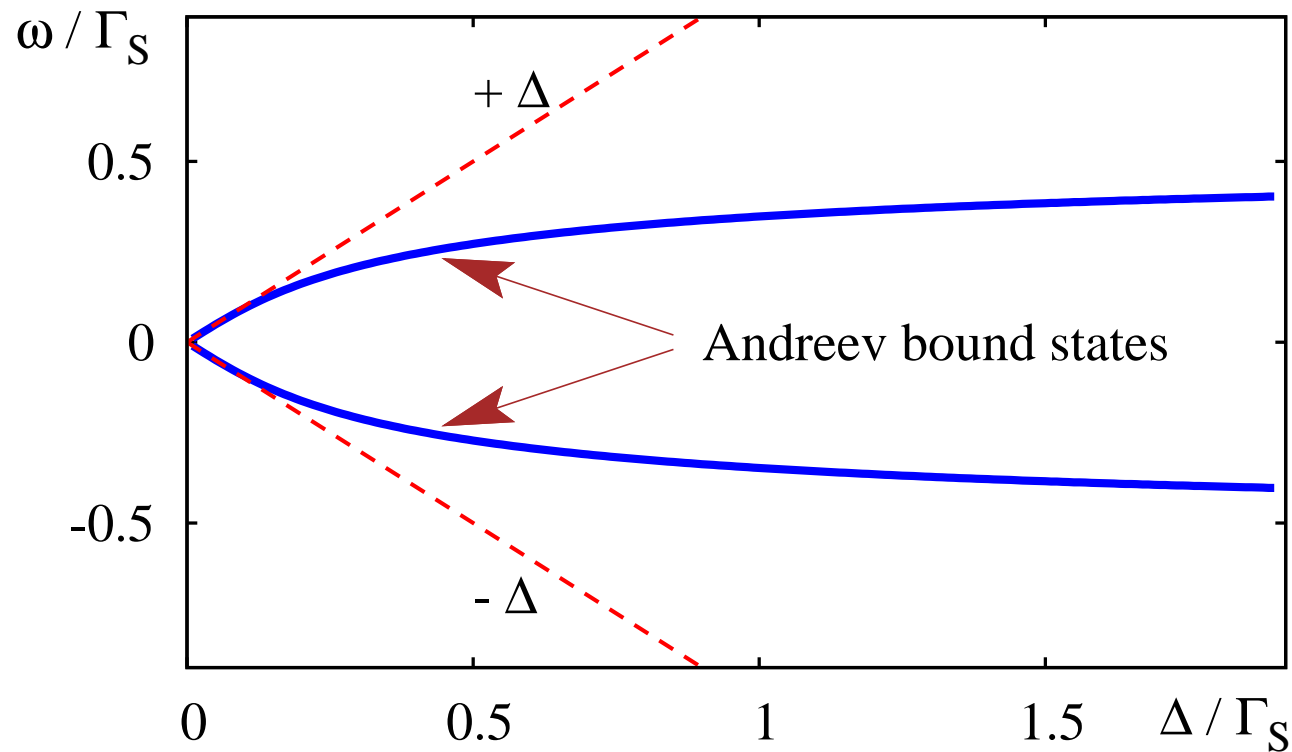
Spectrum consists of:

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Microscopic description

– exact solution for $U = 0$

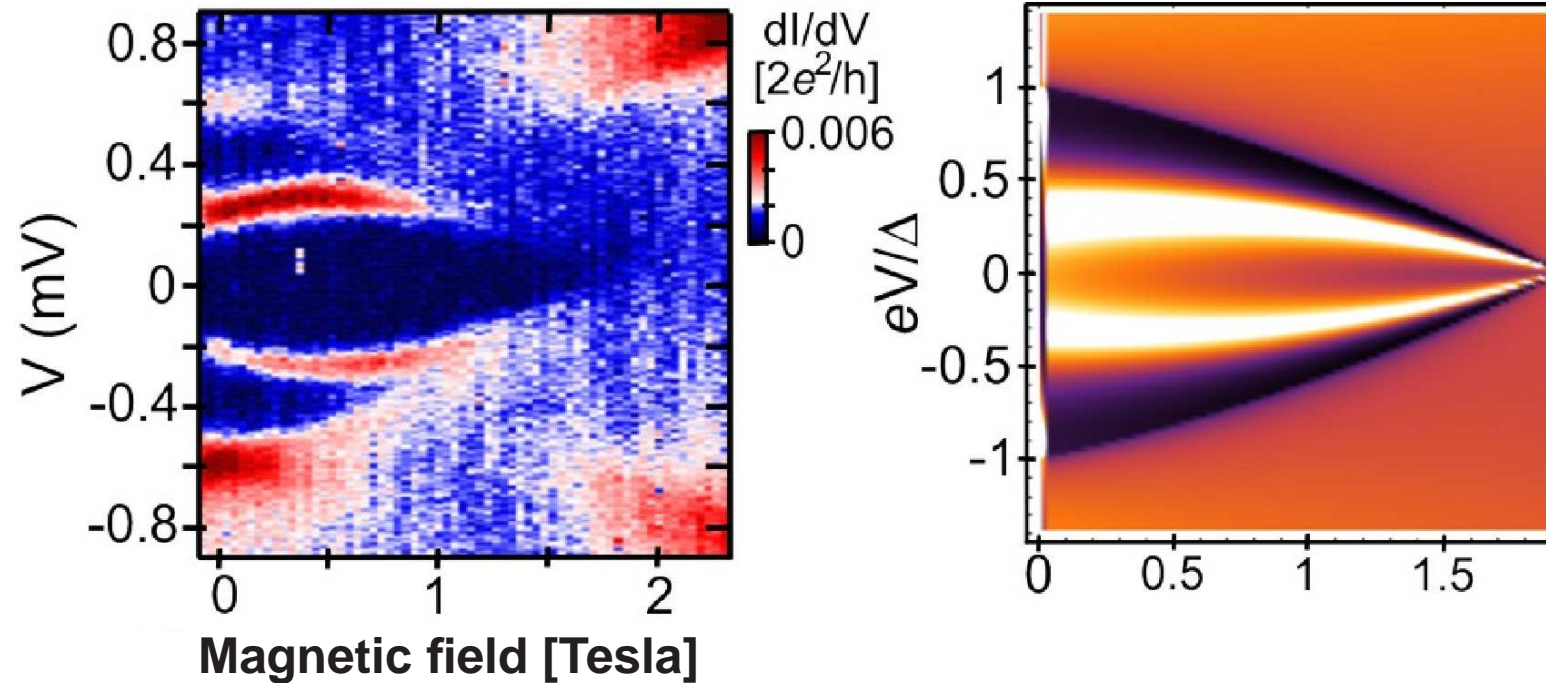


Energies of the in-gap resonances (Andreev bound states)

J. Phys.: Condens. Matter **25**, 435305 (2013).

Microscopic description

– exact solution for $U = 0$



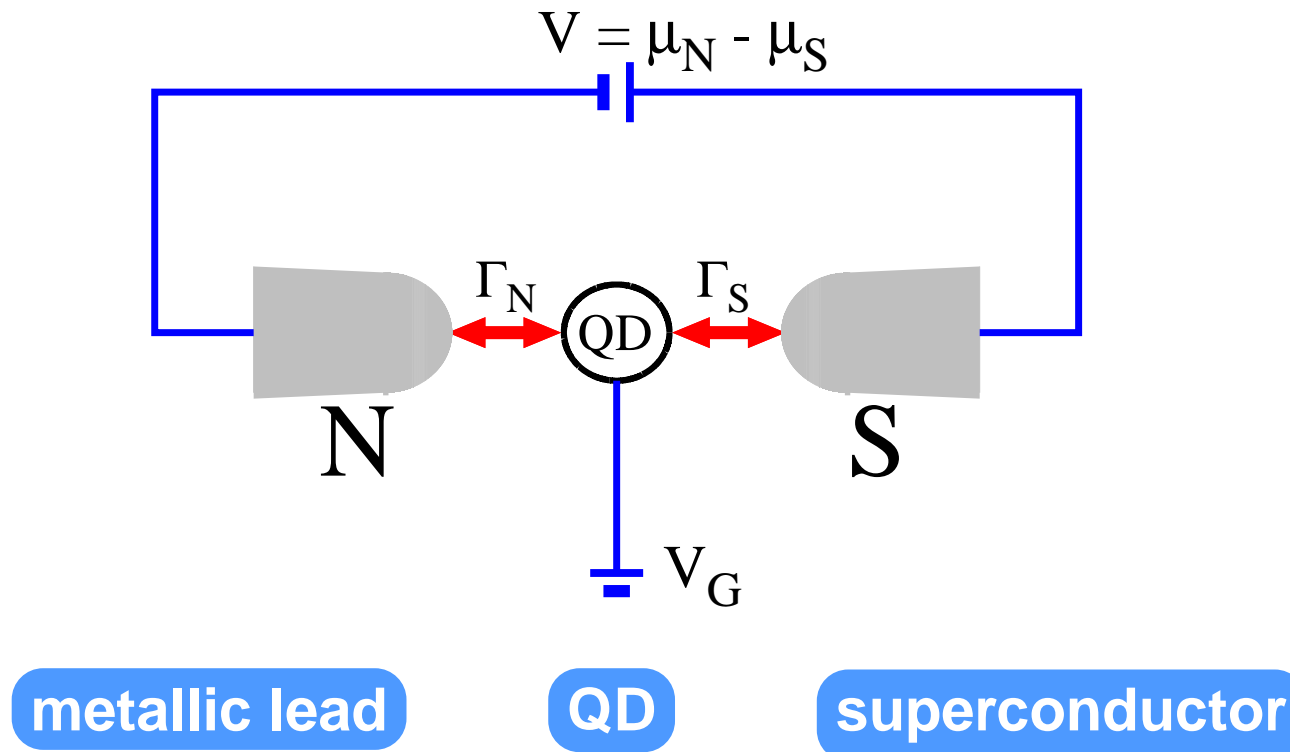
Differential conductance of nanotubes coupled to vanadium (S) and gold (N)
/ external magnetic field changes a magnitude of the pairing gap $\Delta(B)$ /

Eduardo J.H. Lee, ..., S. De Franceschi, Nature Nanotechnology **9**, 79 (2014).

Spectroscopic tools

– probing nano-superconductors

To probe the in-gap states one can study the electron transport through a quantum dot (QD) coupled between the normal (N) and superconducting (S) electrodes



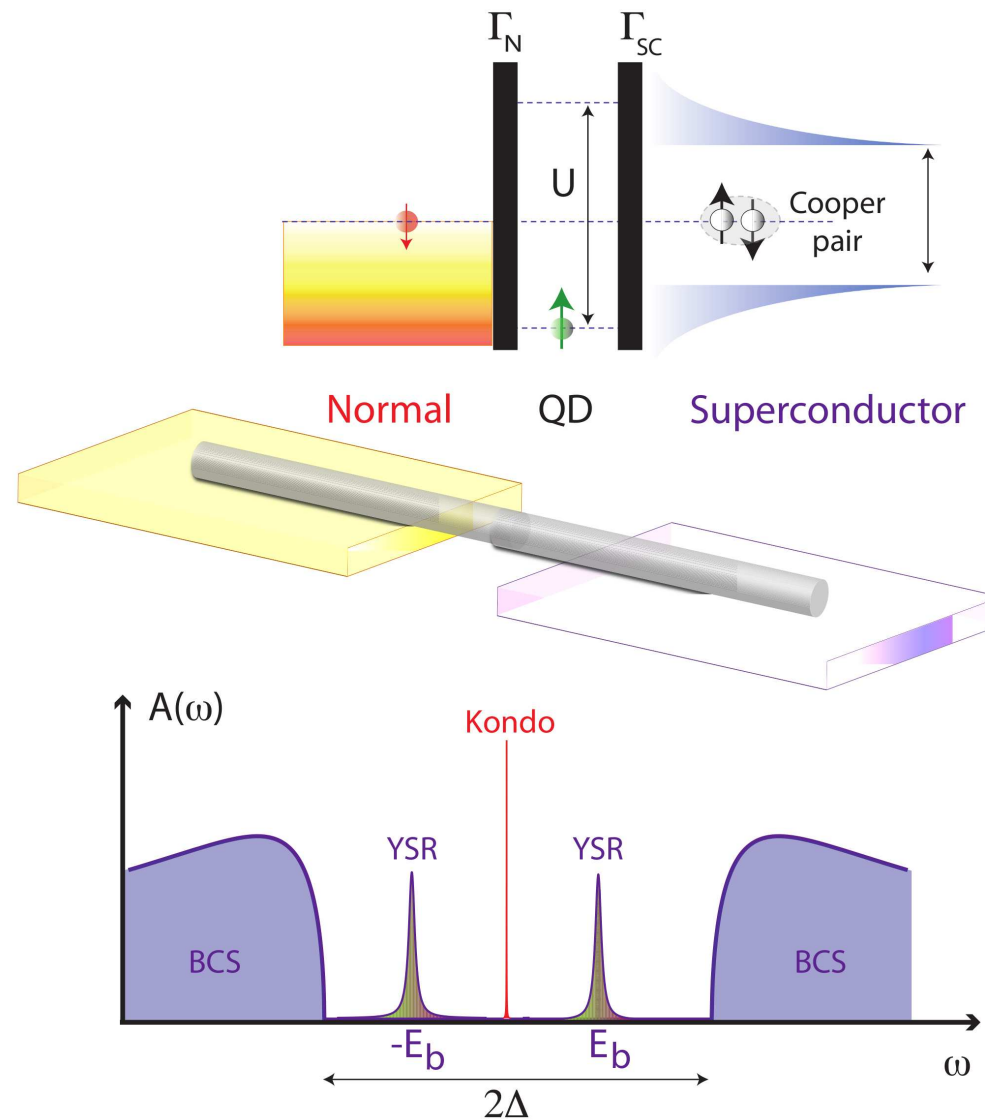
This N-QD-S setup has been practically studied in several experiments.

N - QD - S junctions

– **pairing vs Coulomb repulsion**

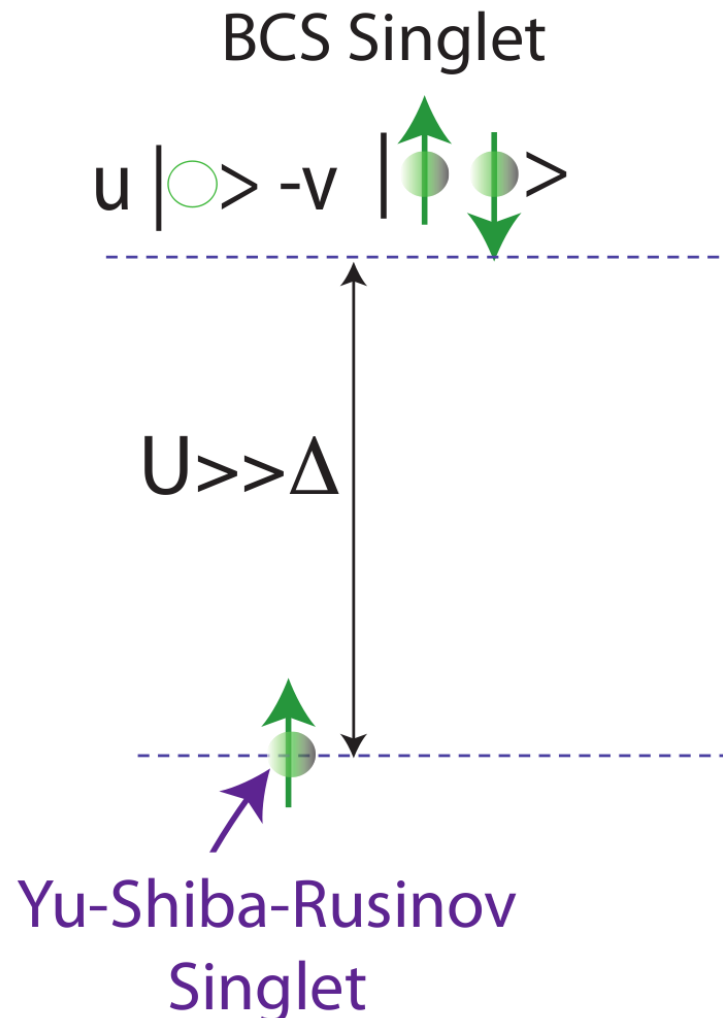
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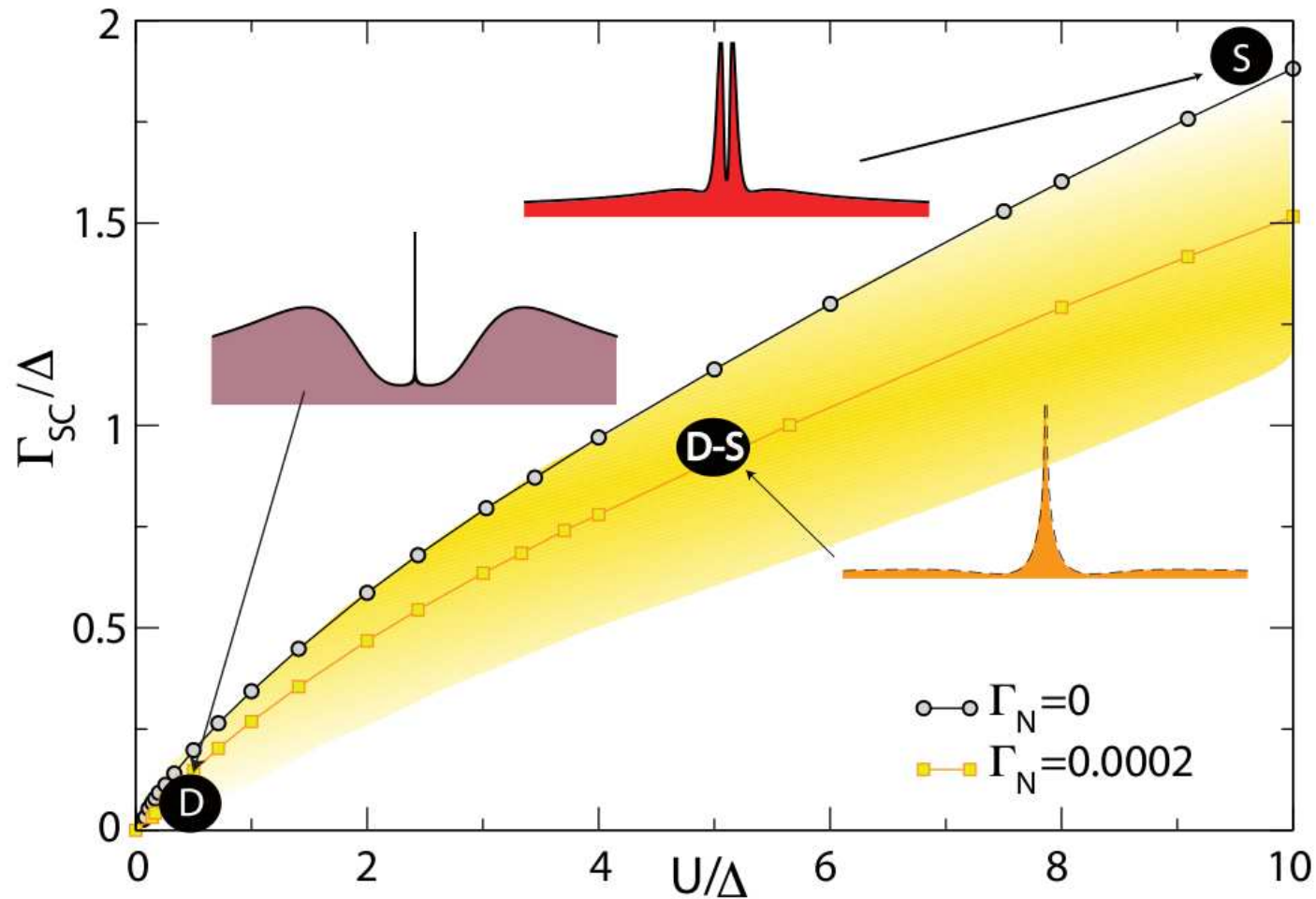
N - QD - S junctions

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N - QD - S junctions

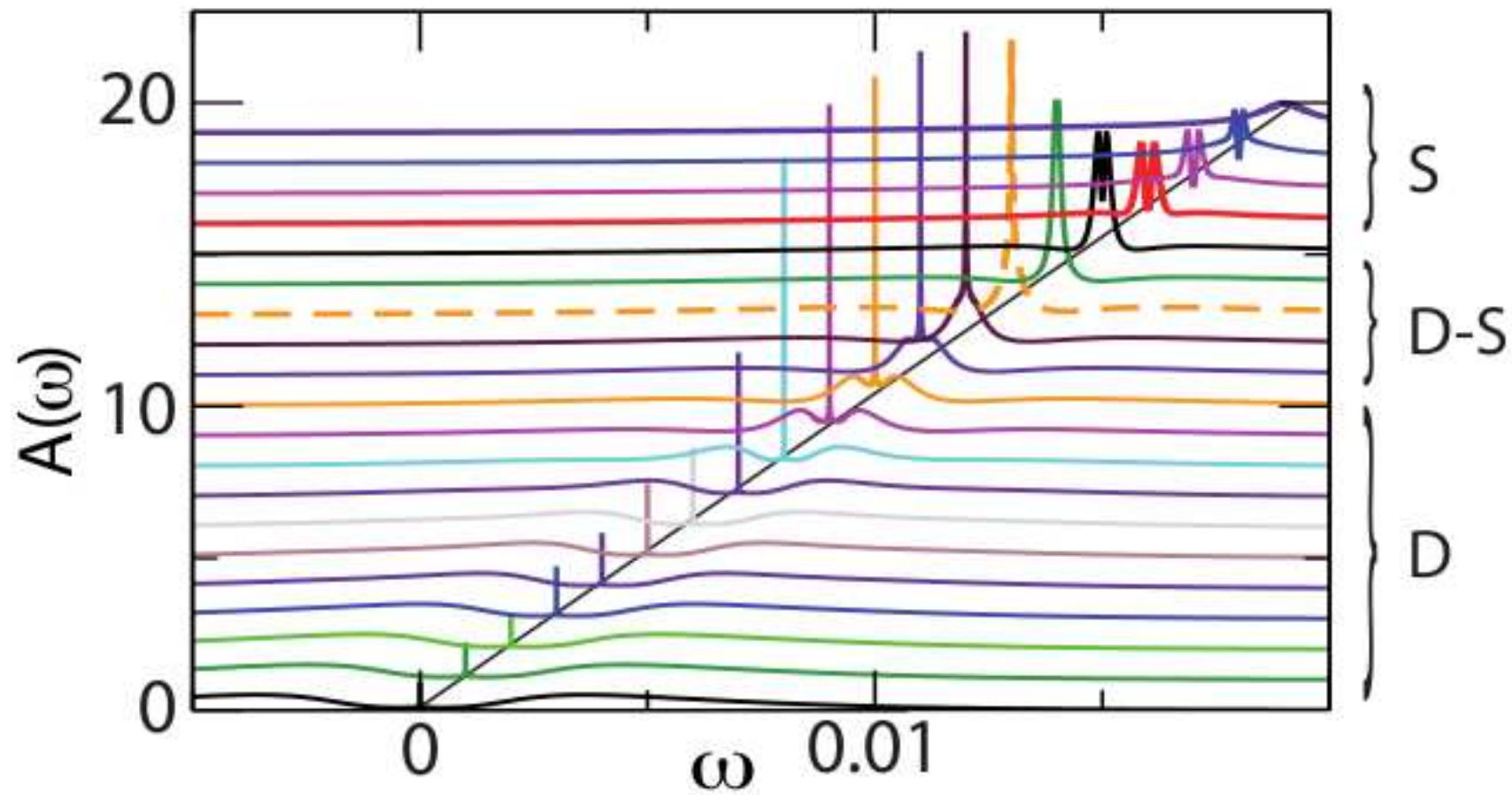
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R. Žitko, J.S. Lim, R. López, and R. Aguado, Phys. Rev. B **91**, 045441 (2015).

N - QD - S junctions

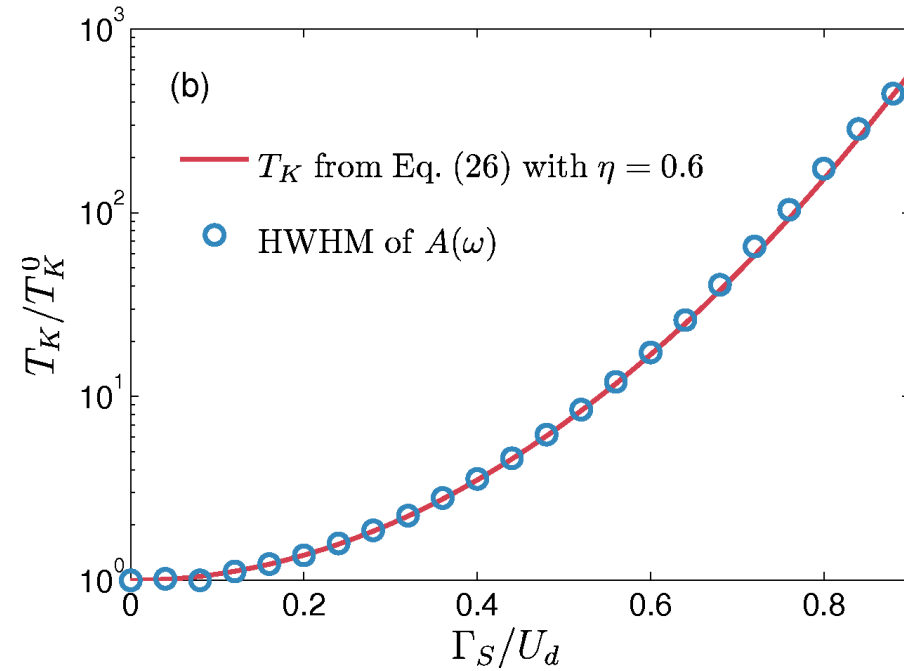
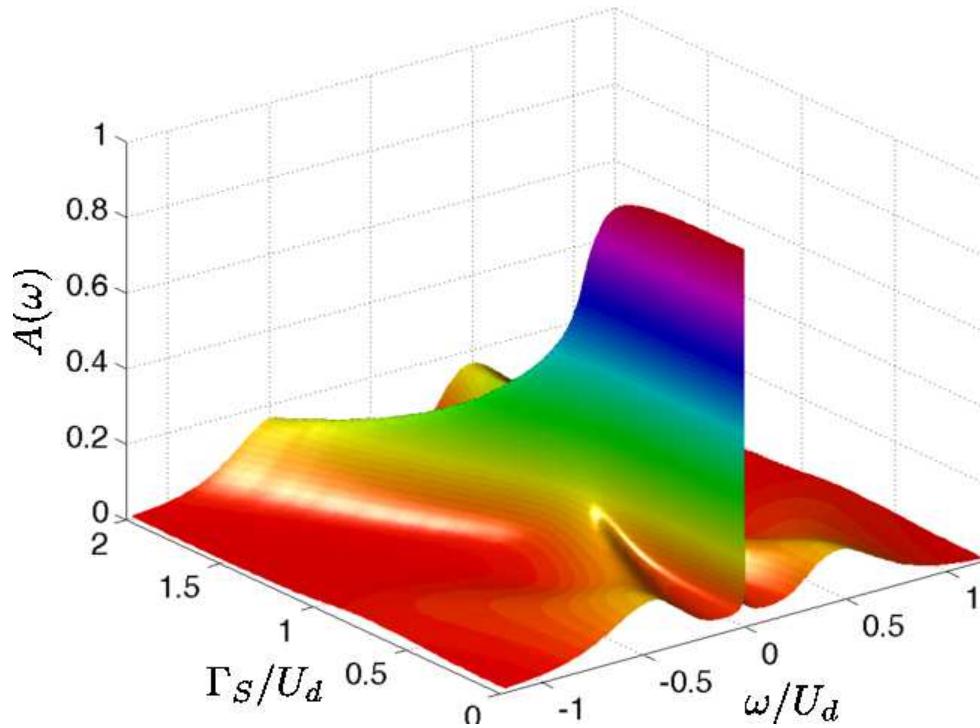
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N - QD - S junctions

– pairing vs Coulomb repulsion



T. Domański, I. Weymann & M. Barańska, arXiv:1507.01851 (2015) preprint.

Our and R. Žitko's studies reveal that: T_K is enhanced by Γ_S

3. Nano-superconductivity:



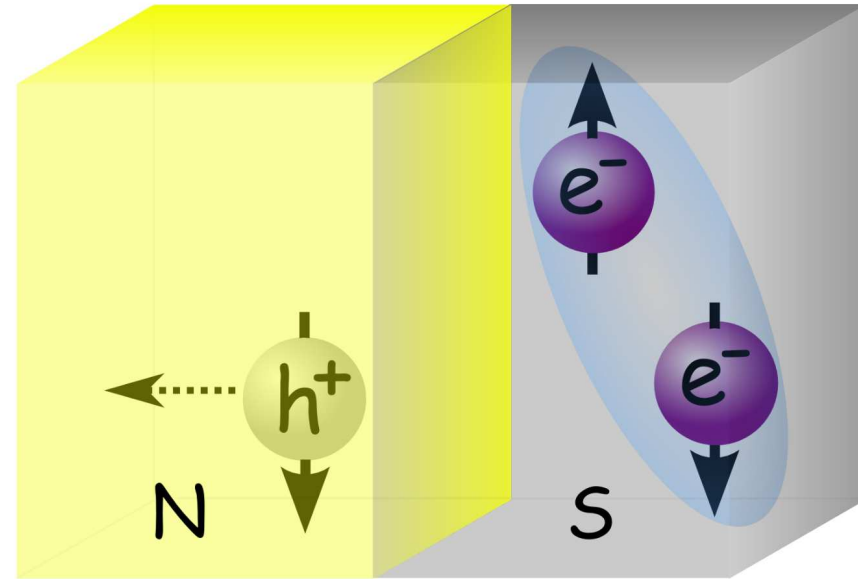
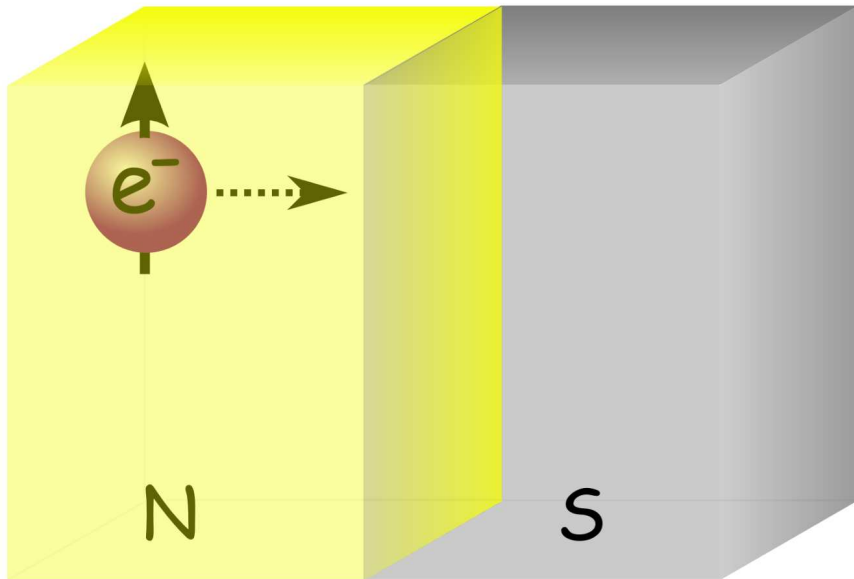
some practical aspects

Andreev reflections

– **possible applications**

Andreev reflections

– possible applications

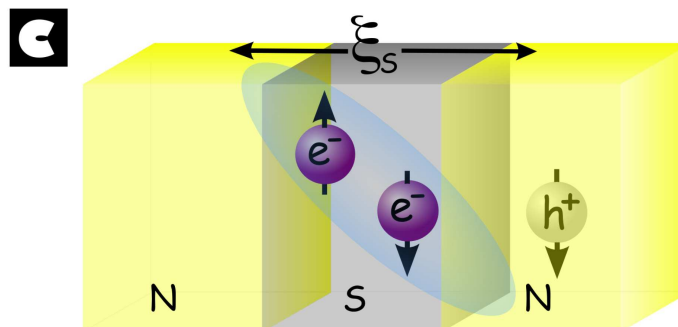
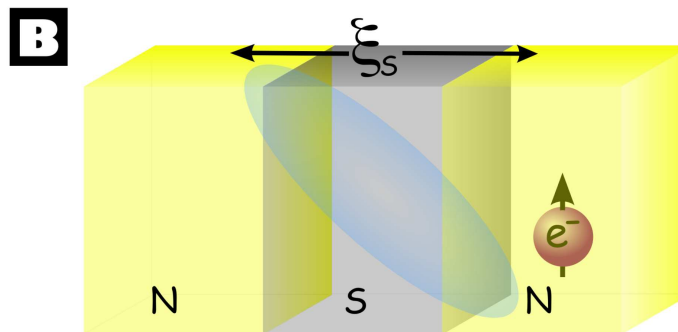
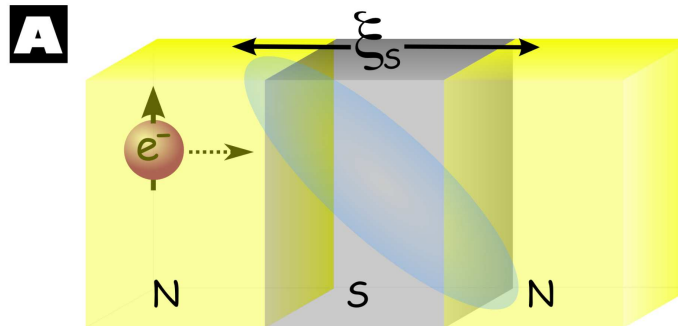


Schematic illustration of the Andreev-type scattering

Andreev reflections

– possible applications

Andreev-type scattering can be also considered in more complex junctions

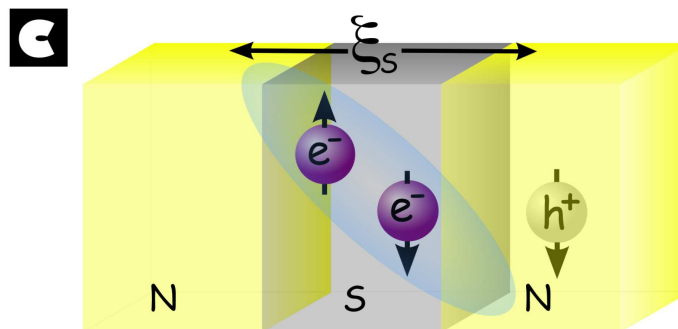
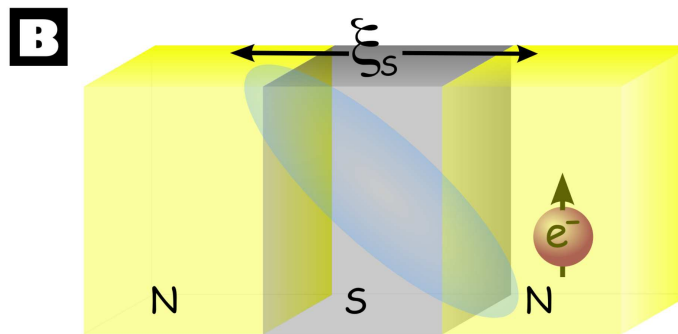
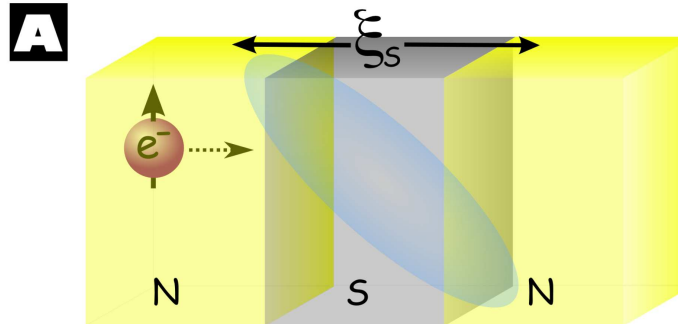


Andreev reflections

– possible applications

Andreev-type scattering can be also considered in more complex junctions

incident electron

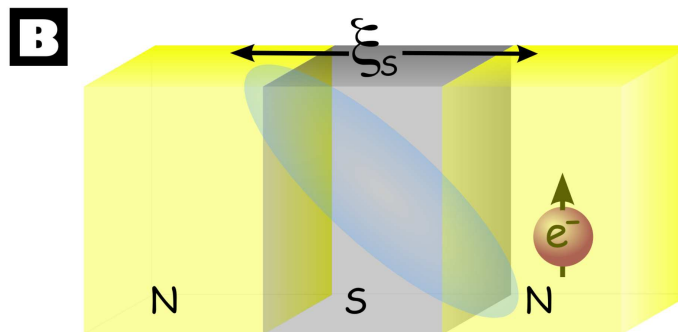
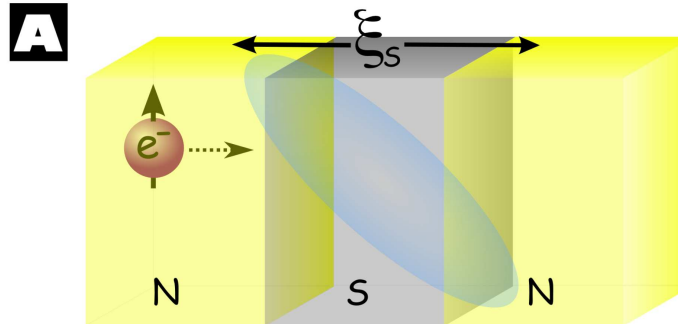


Andreev reflections

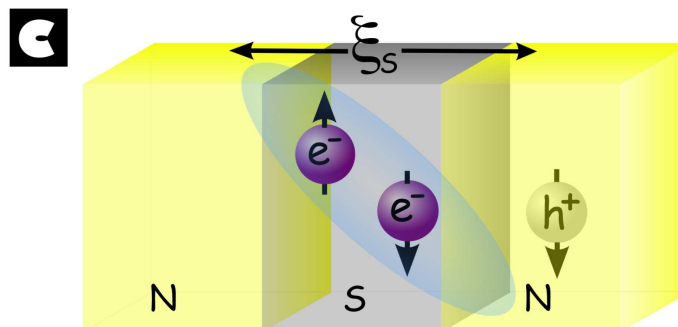
possible applications

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elastic tunnelling

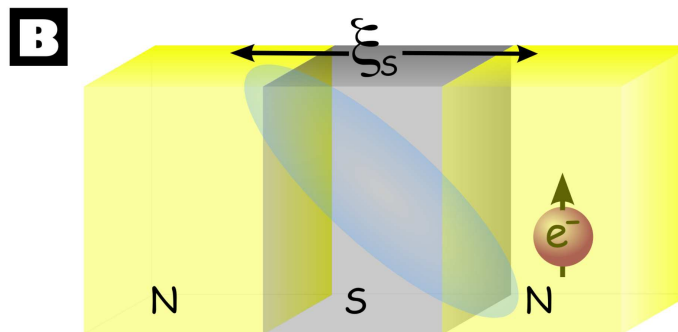
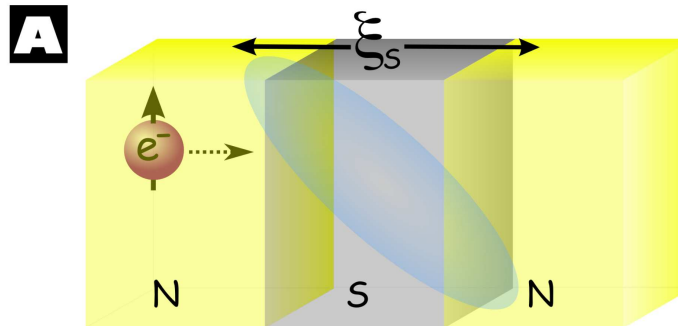


Andreev reflections

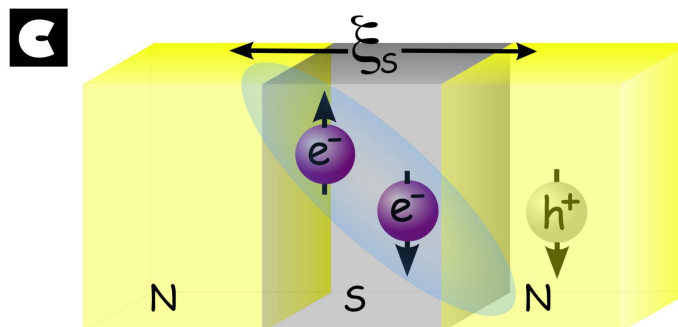
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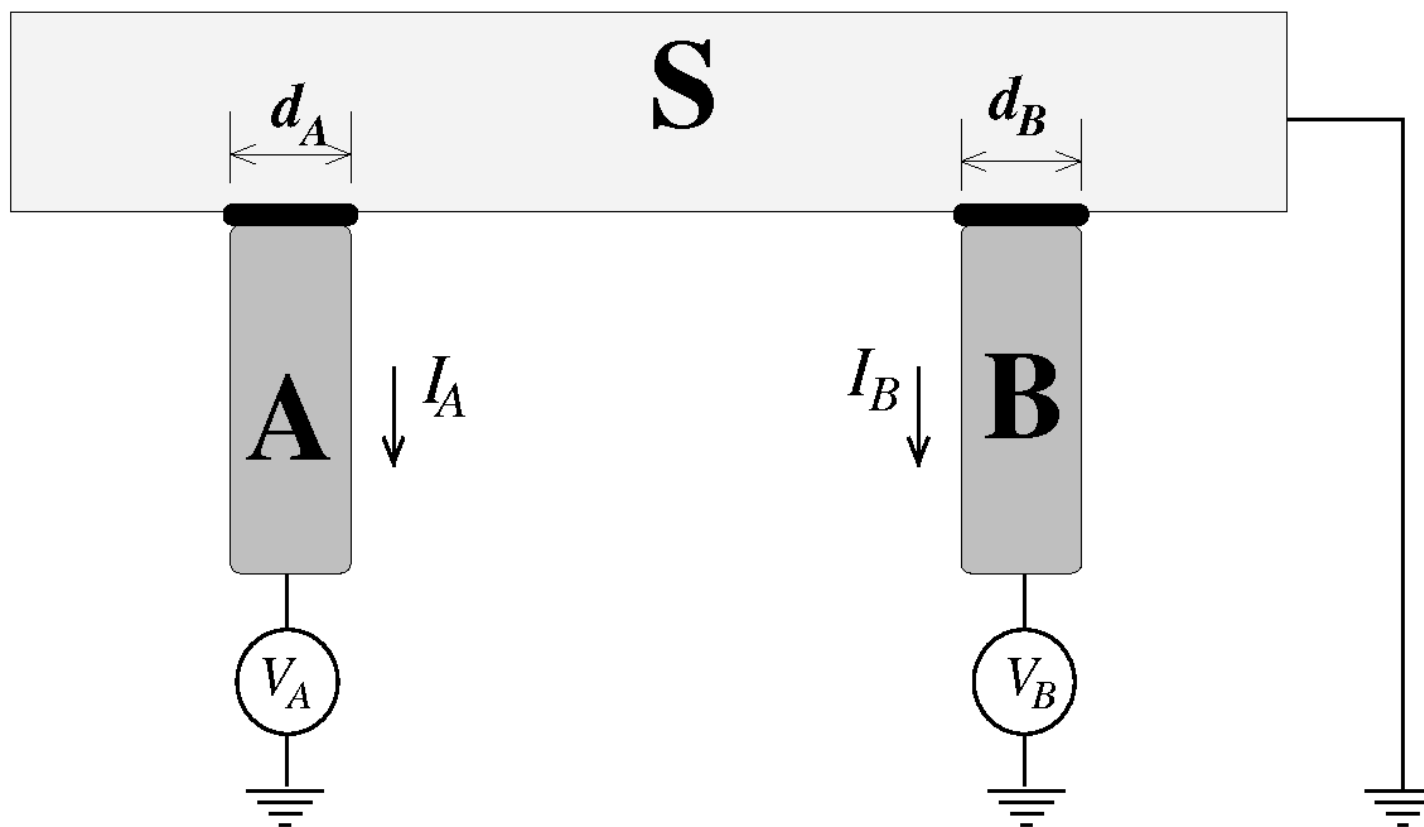


crossed Andreev refl.

Non-local transport – planar junctions

Non-local transport – planar junctions

These ET/CAR processes have first considered in the planar junctions

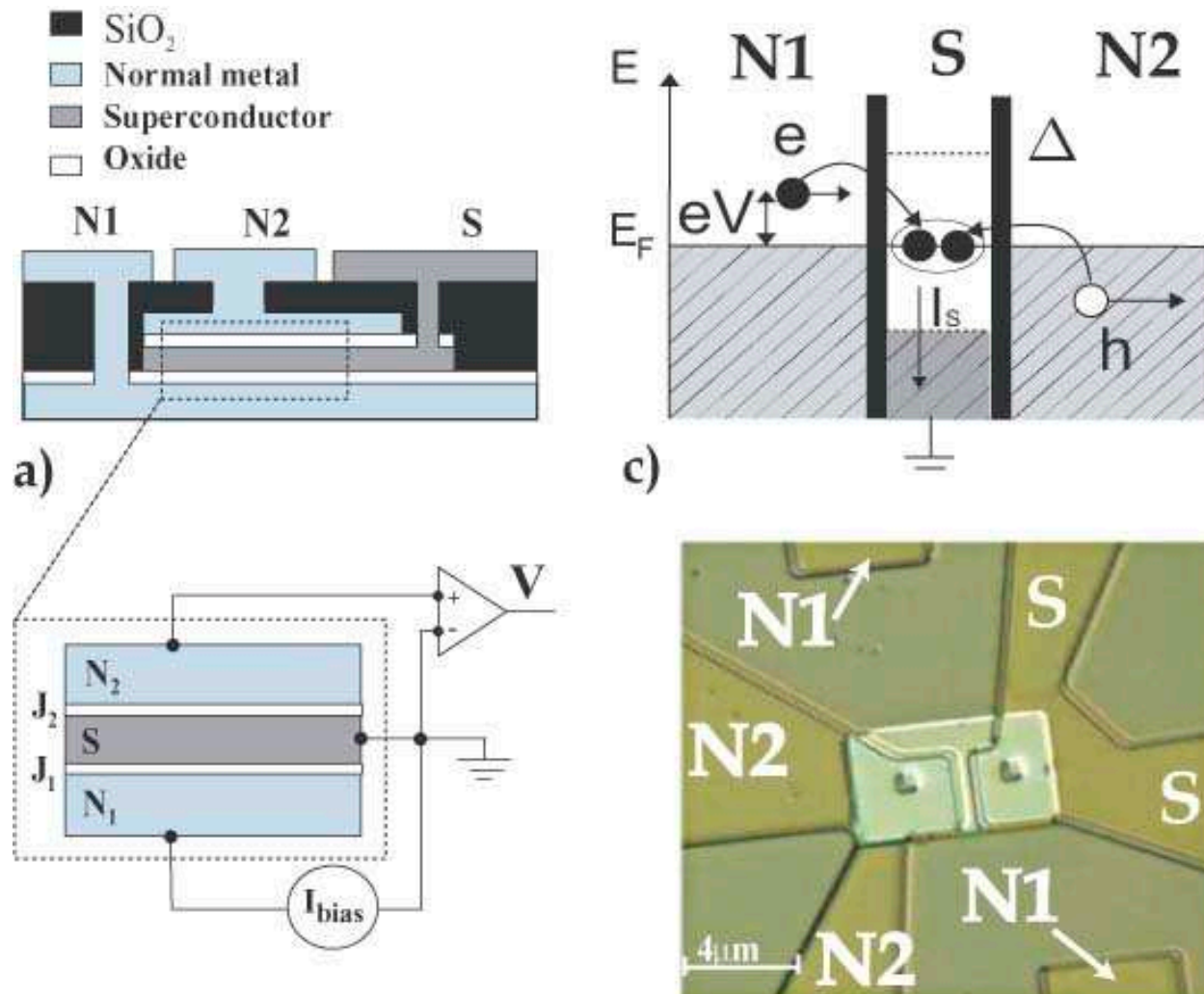


A, B – normal electrodes, S – superconducting material

G. Falci, D. Feinberg, F. Hekking, Europhys. Lett. **54**, 255 (2001).

Non-local transport – planar junctions

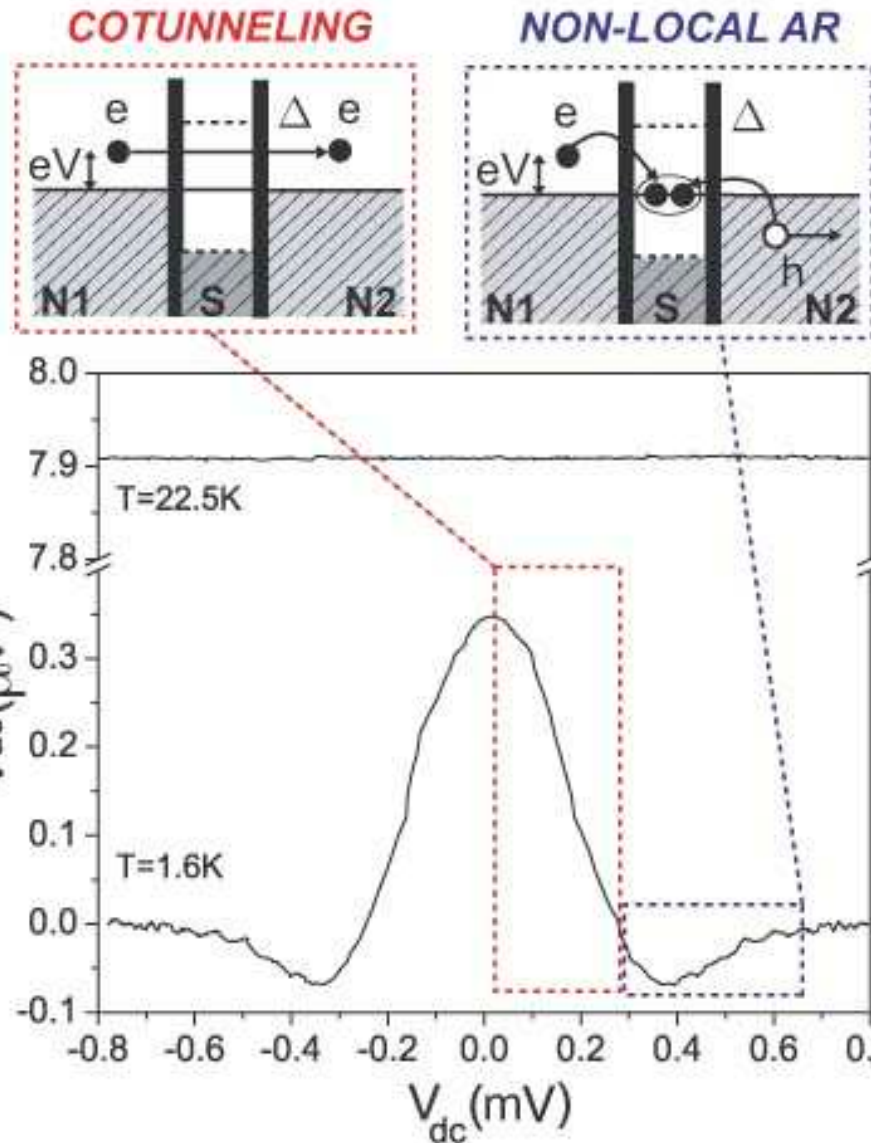
Experimental realization (Delft group)



S. Russo, M. Kroug, T. M. Klapwijk & A.F. Morpurgo, *Phys. Rev. Lett.* **95**, 027002 (2005).

Non-local transport – planar junctions

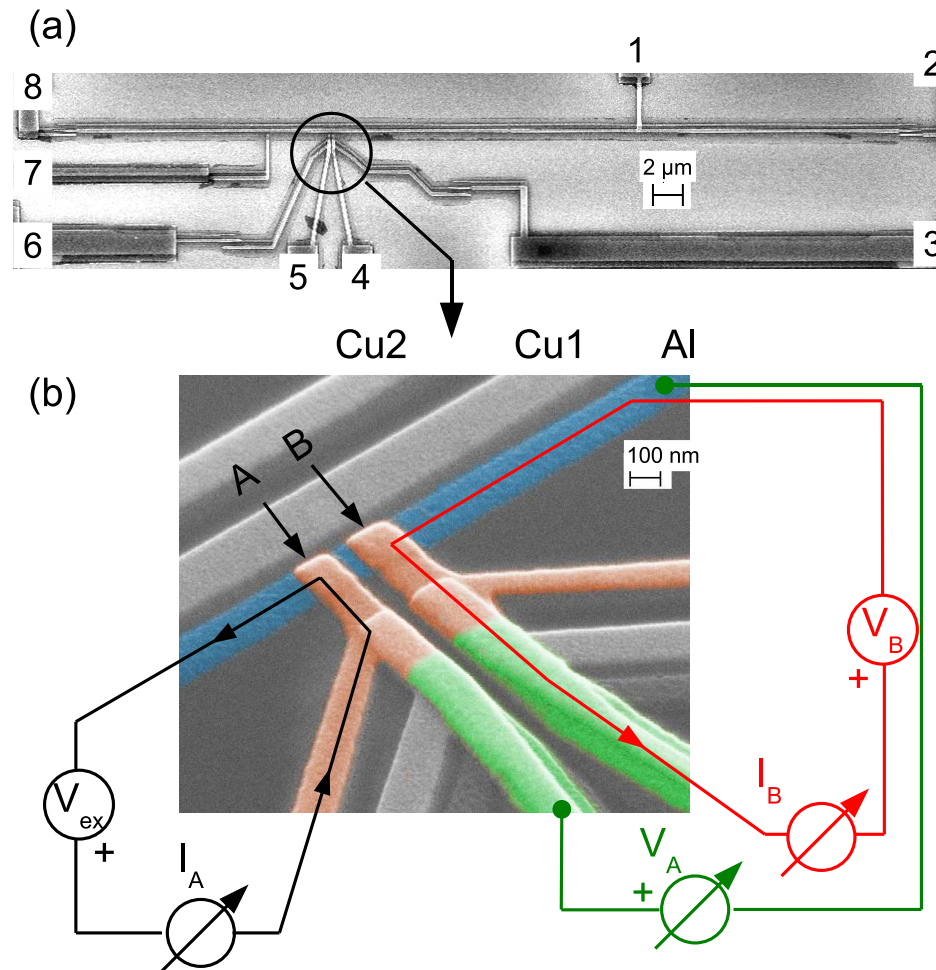
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Non-local transport – planar junctions

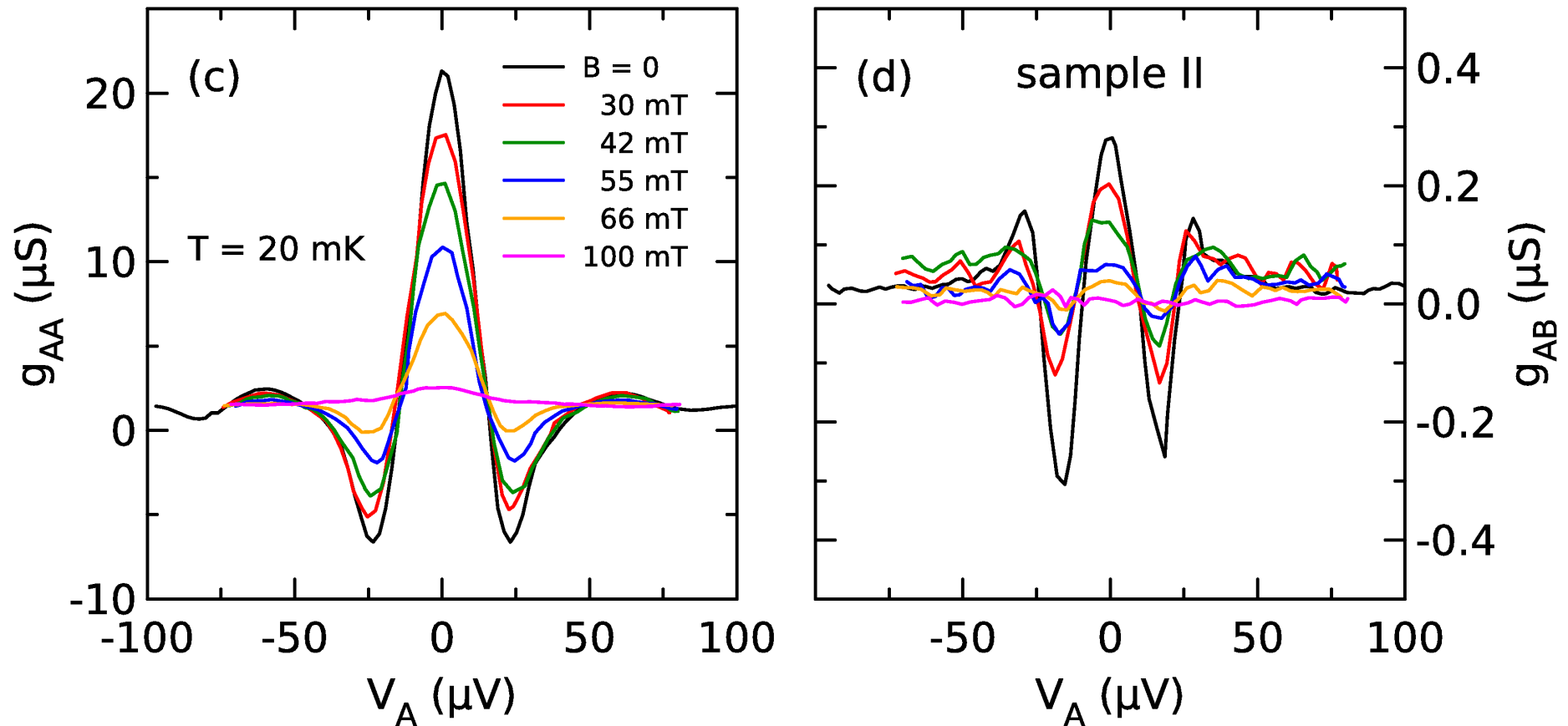
Experimental realization (Karlsruhe group)



J. Brauer, F. Hübler, M. Smetanin, D. Beckman, D. & H. von Löhneysen, *Phys. Rev. B* **81**, 024515 (2010).

Non-local transport – planar junctions

Experimental realization (Karlsruhe group)



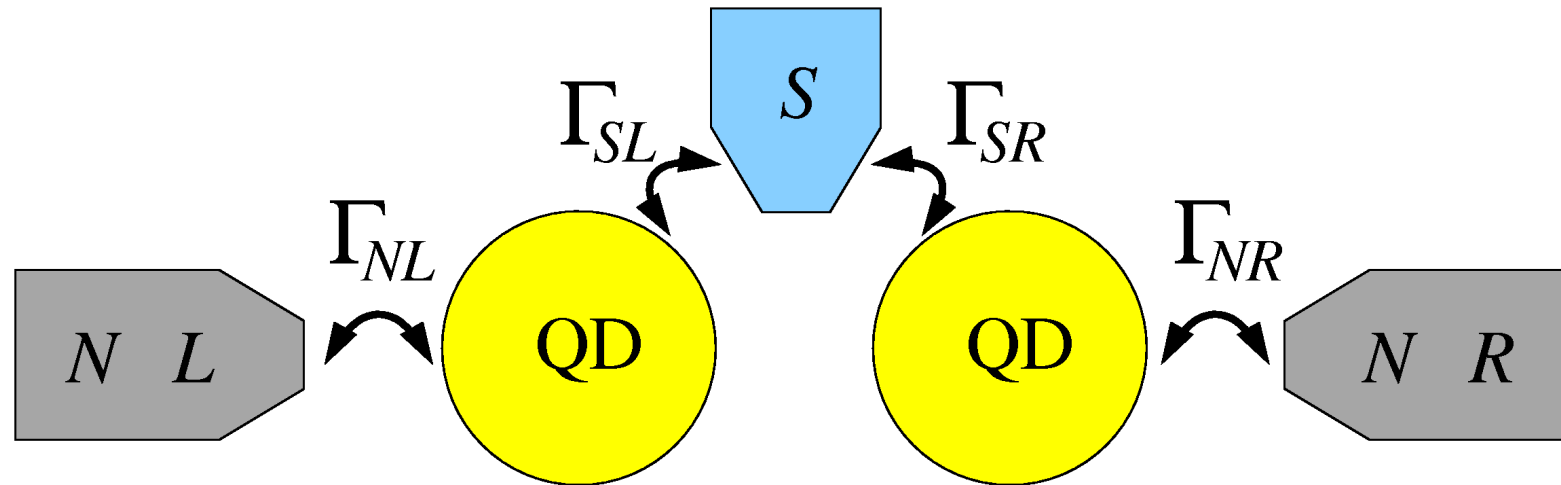
J. Brauer, F. Hübler, M. Smetanin, D. Beckman, D. & H. von Löhneysen, *Phys. Rev. B* **81**, 024515 (2010).

3-terminal junctions

– with quantum dots

3-terminal junctions

– with quantum dots



Cooper pairs are split, preserving entanglement of individual electrons.

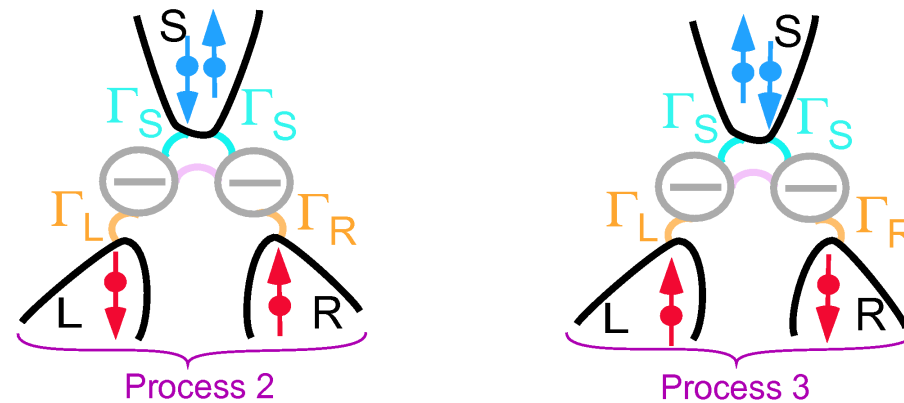
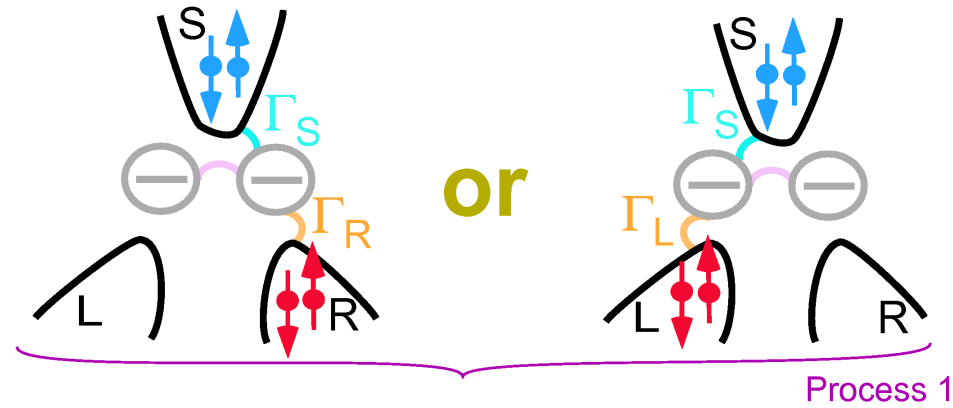
L. Hofstetter, S. Csonka, J. Nygård, C. Schönenberger, Nature **461**, 960 (2009).

J. Schindele, A. Baumgartner, C. Schönenberger, Phys. Rev. Lett. **109**, 157002 (2012).

... and many other groups.

3-terminal junctions

– with quantum dots

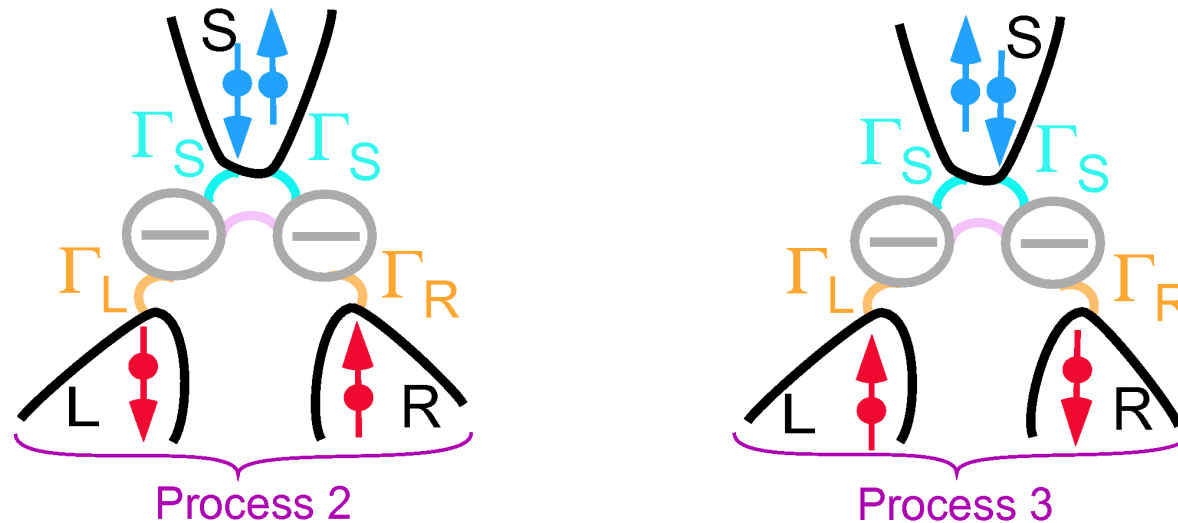


Possible channels of the Cooper pair splitting

L.G. Herrmann et al, Phys. Rev. Lett. 104, 026801 (2010).

3-terminal junctions

– with quantum dots



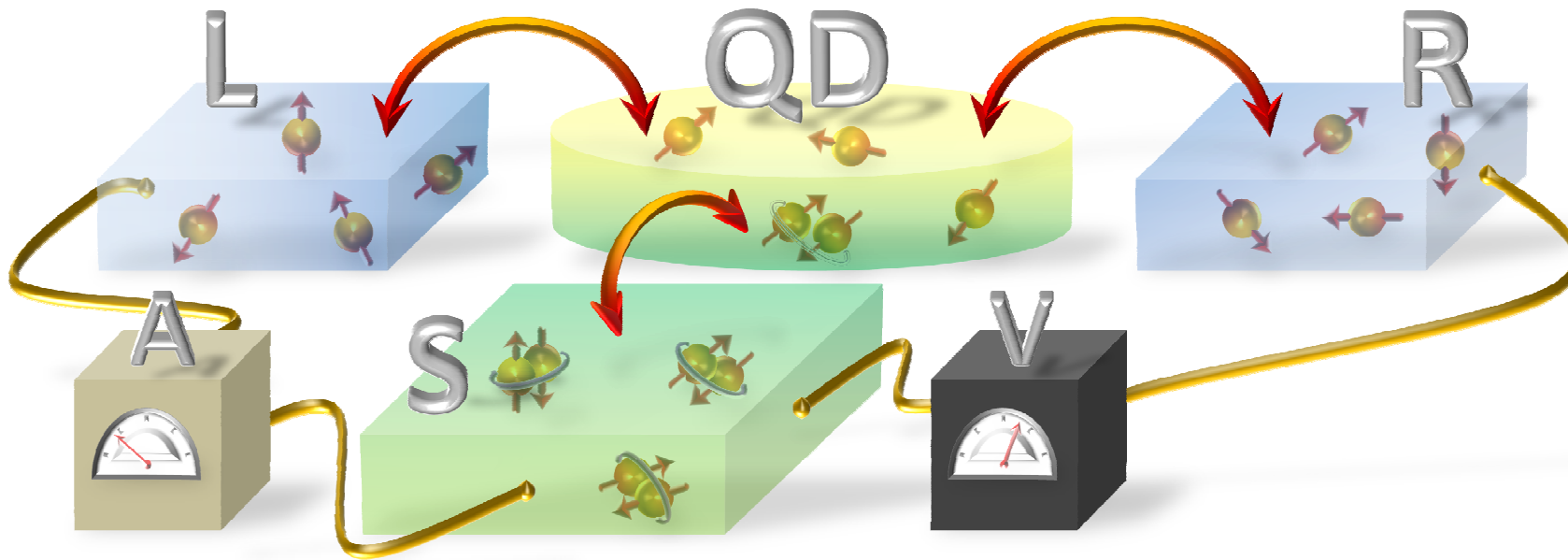
These processes are similar to the crossed Andreev scattering and cause the strong non-local transport properties.

Non-local transport

– **crossed Andreev reflections**

Non-local transport – crossed Andreev reflections

Quantum impurity in the 3-terminal configuration

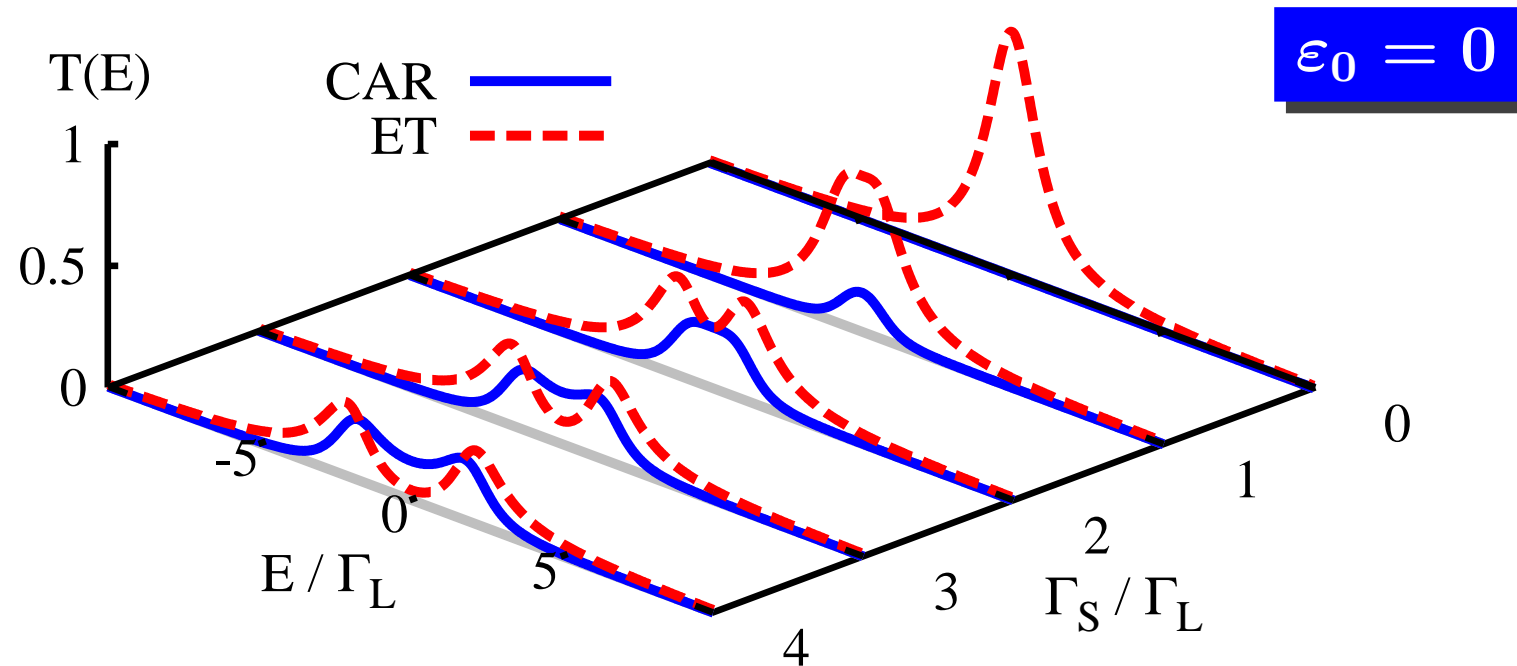


L, R – normal electrodes, S – superconducting reservoir, QD – quantum dot

G. Michałek, T. Domański, B.R. Bułka & K.I. Wysokiński, Scientific Reports 5, 14572 (2015).

Non-local transport – crossed Andreev reflections

Transmittance of the non-local transport channels

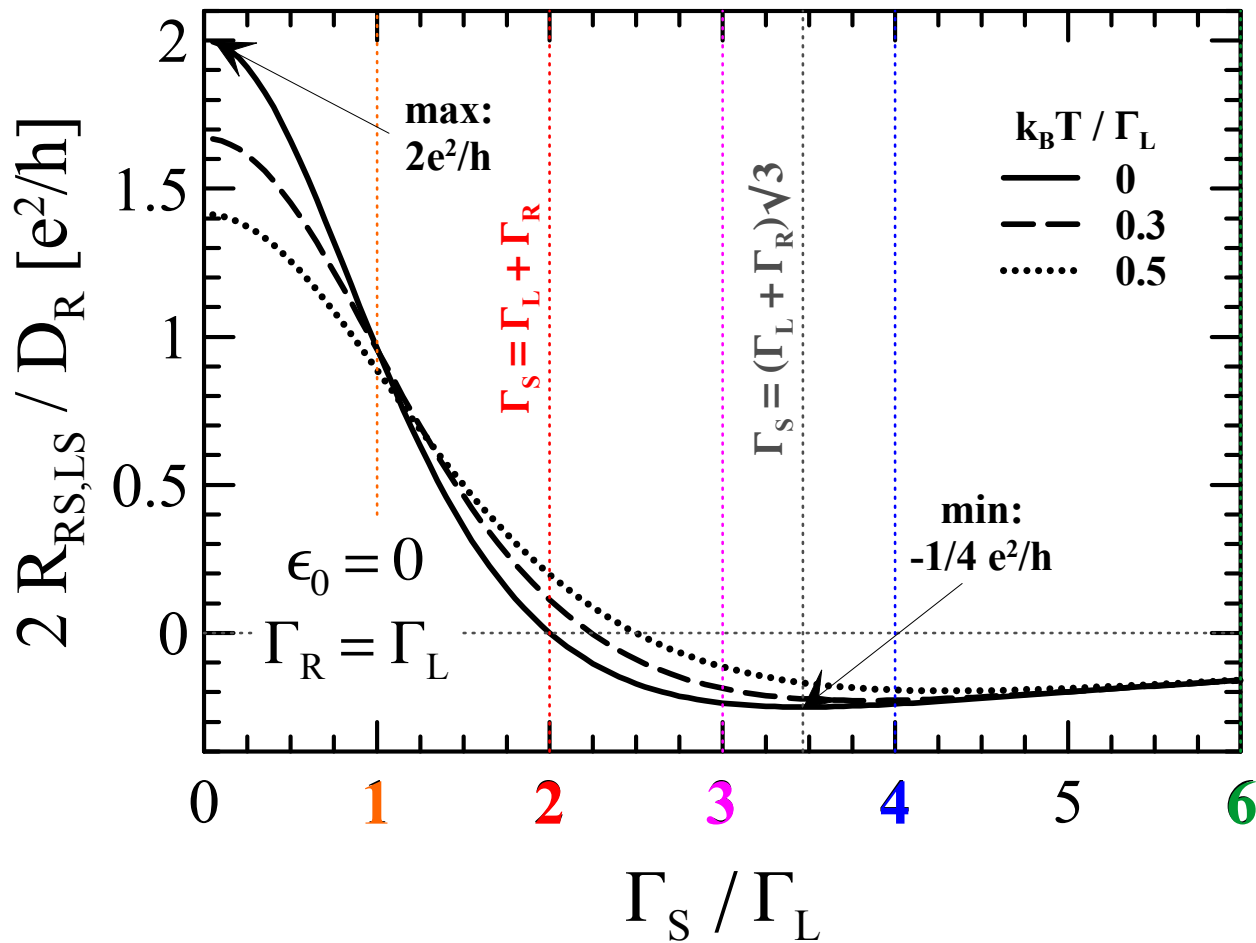


ET – single electron transfer, **CAR** – crossed Andreev reflection

Non-local transport

– crossed Andreev reflections

Non-local resistance in the linear response limit.

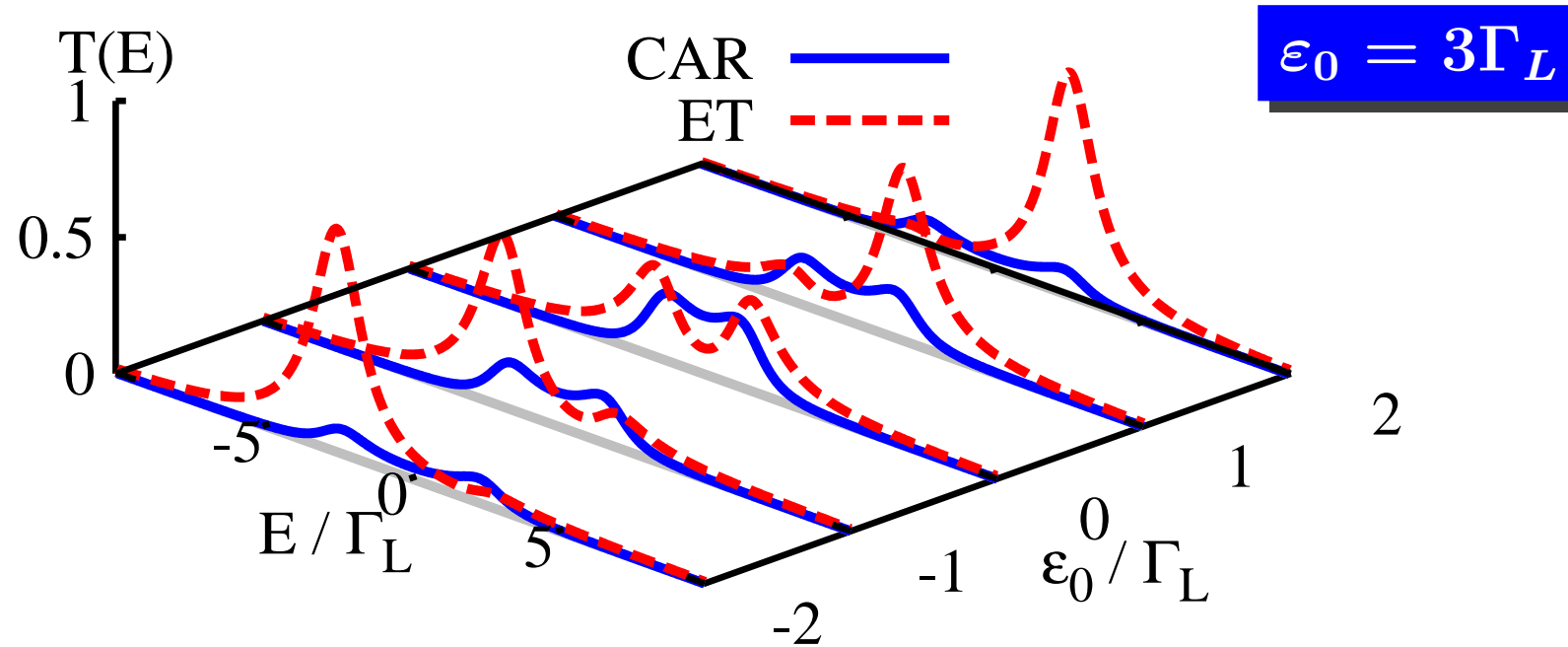


$$R_{RS,LS} \equiv V_{RS} / I_{LS}$$

Negative resistance for strong enough coupling Γ_S !

Non-local transport – crossed Andreev reflections

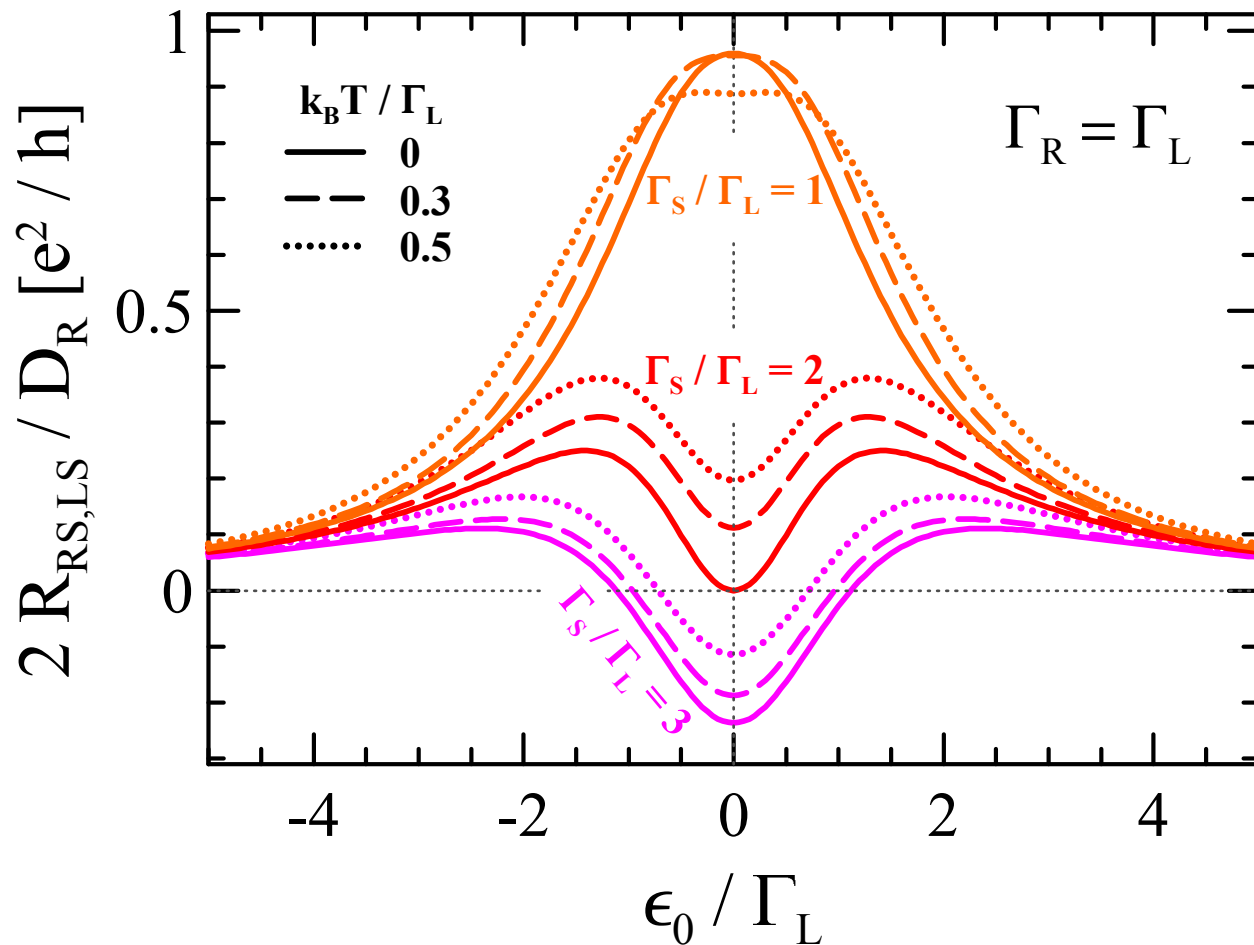
Transmittance of the non-local transport channels



ET – single electron transfer, **CAR** – crossed Andreev reflection

Non-local transport – **crossed Andreev reflections**

Non-local resistance in the linear response limit.



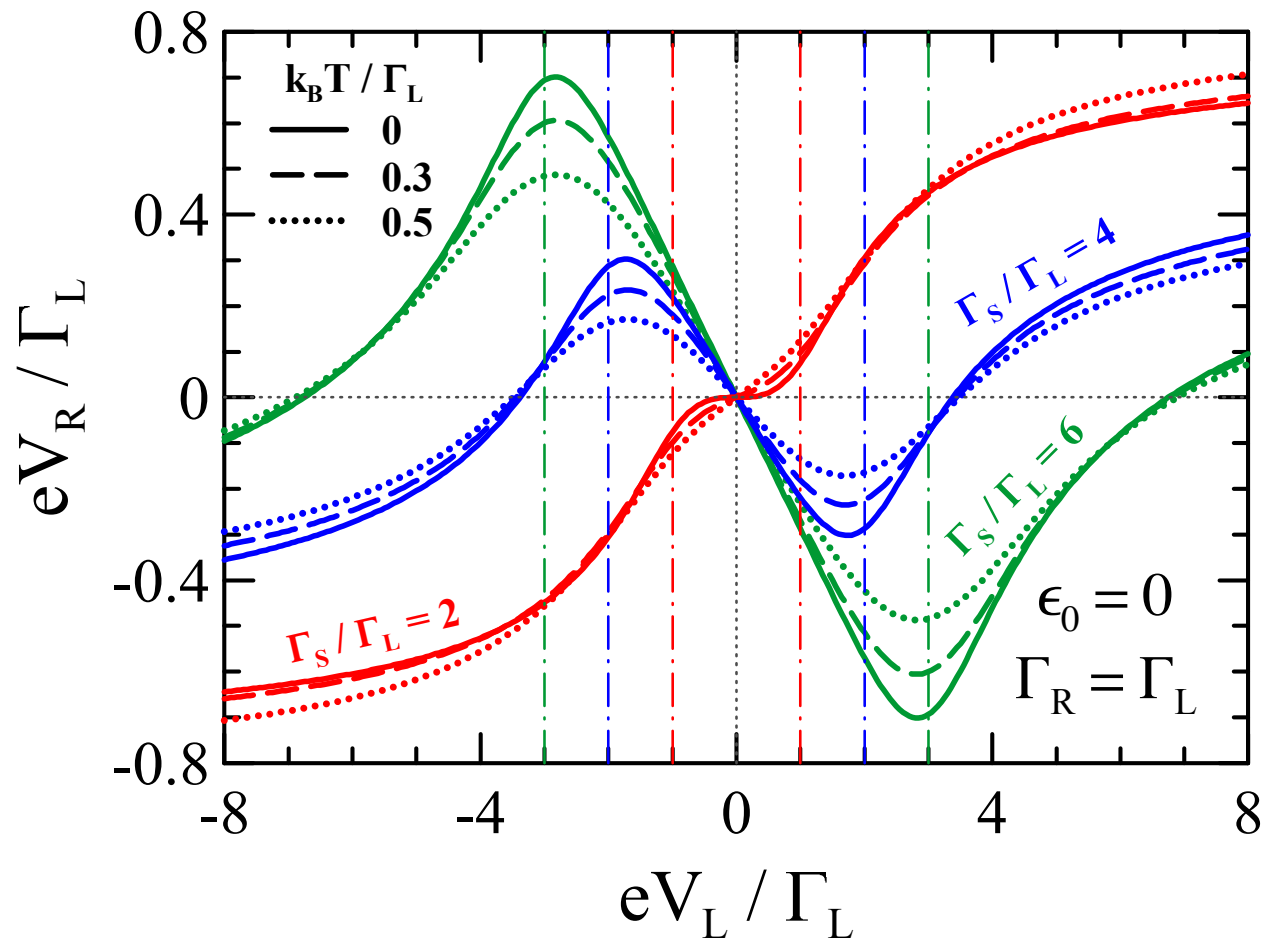
$$R_{RS,LS} \equiv V_{RS} / I_{LS}$$

Negative resistance for strong enough coupling Γ_S and $\epsilon_0 \sim 0$!

Non-local transport

– crossed Andreev reflections

Beyond the linear response regime.



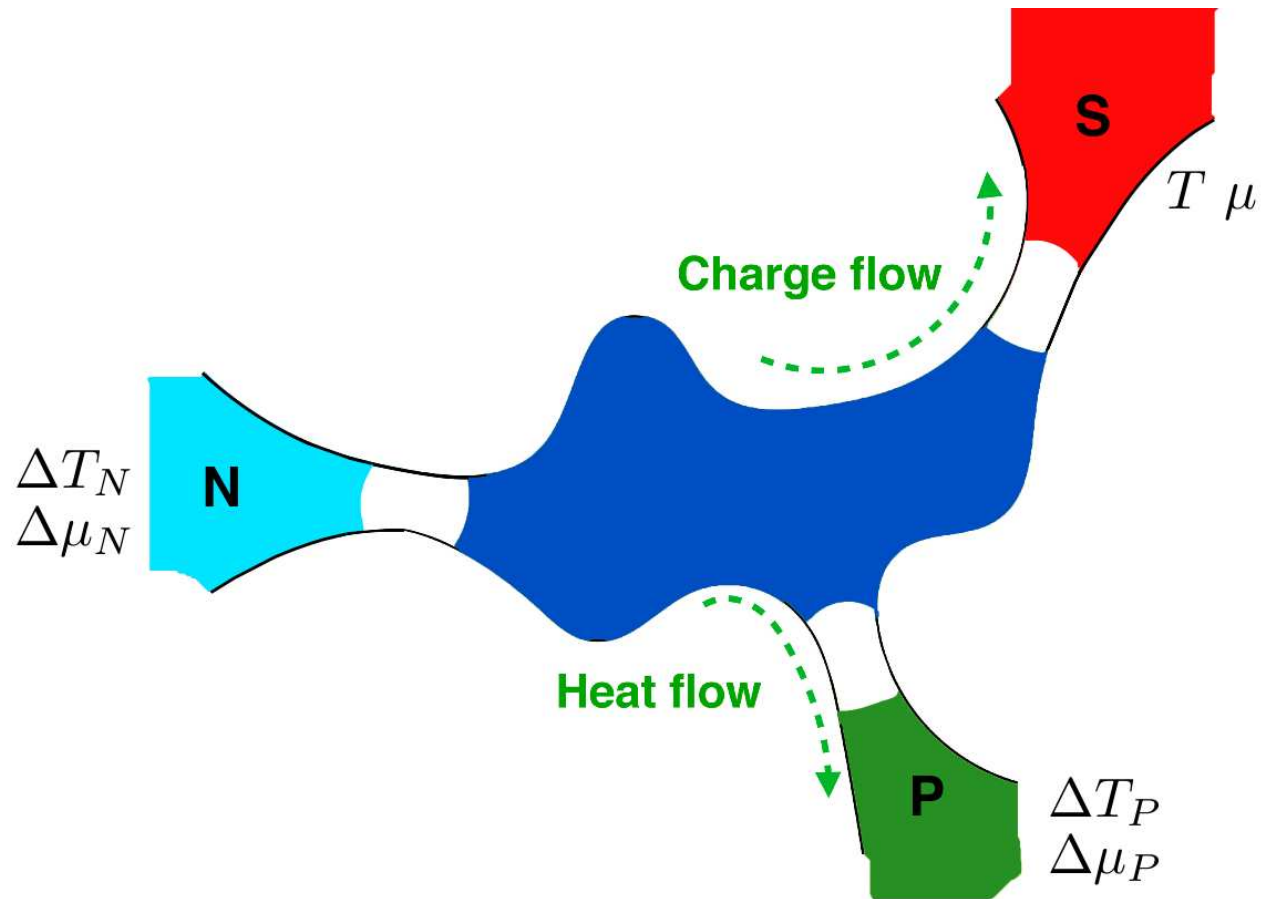
Inverse sign of the non-local voltage for strong enough coupling Γ_S .

Andreev reflections

– **other perspectives**

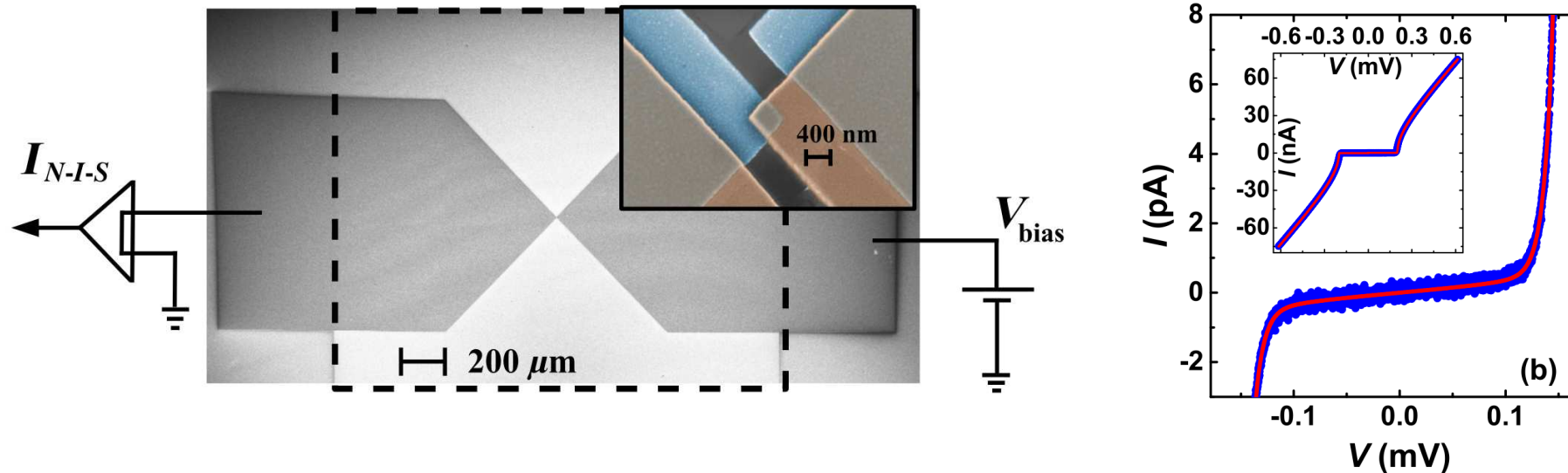
Andreev reflections – other perspectives

Crossed Andreev reflections enable the separation of charge from heat currents



Andreev reflections – other perspectives

On-chip nanoscopic thermometer operating down to 7 mK.



A.V. Feshchenko, L. Casparis, ..., J.P. Pekola & D.M. Zumbühl, Phys. Rev. Appl. **4**, 034001 (2015).

T. Faivre, D.S. Golubev and J.P. Pekola, Appl. Phys. Lett. **106**, 182602 (2015).

- Virtues of a device:
- ⇒ is almost free of any self-heating
 - ⇒ operates at cryogenic temperatures
 - ⇒ can thermally-monitor the qubits

Summary

Summary

Nanoscopic superconductors:

Summary

Nanoscopic superconductors:

⇒ **can be induced by the proximity effect**

Summary

Nanoscopic superconductors:

- ⇒ can be induced by the proximity effect
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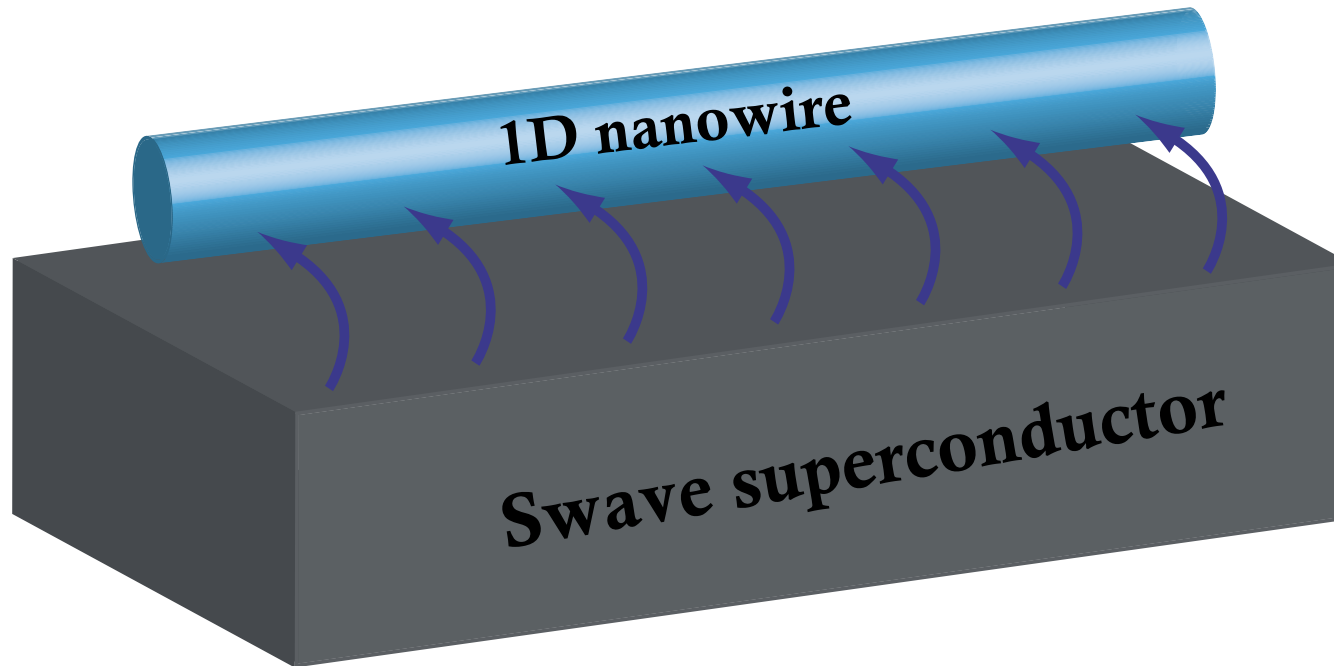
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<http://kft.umcs.lublin.pl/doman/lectures>

Andreev vs Majorana states

– **a story of mutation**

Andreev vs Majorana states – a story of mutation

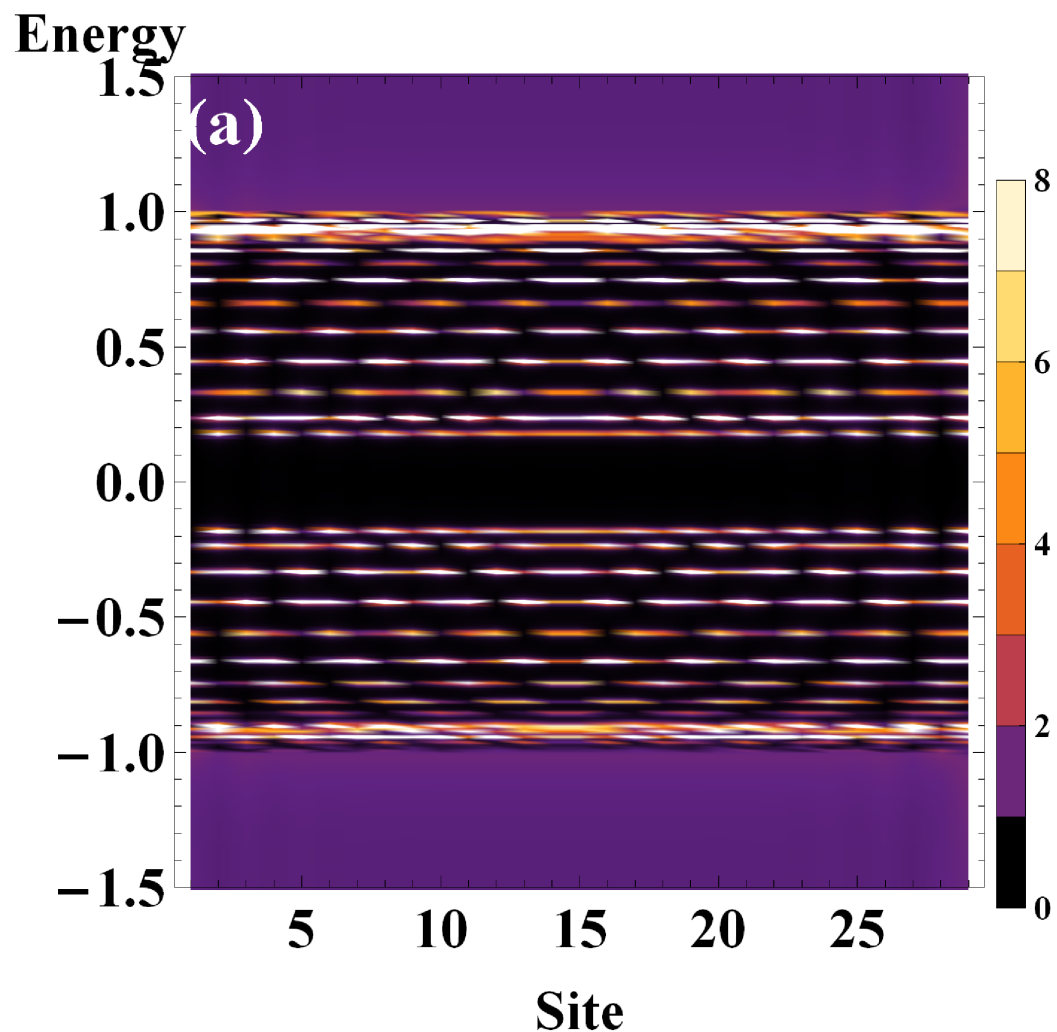


Quantum wire deposited on s-wave superconductor

D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B **88**, 165401 (2013).

Andreev vs Majorana states

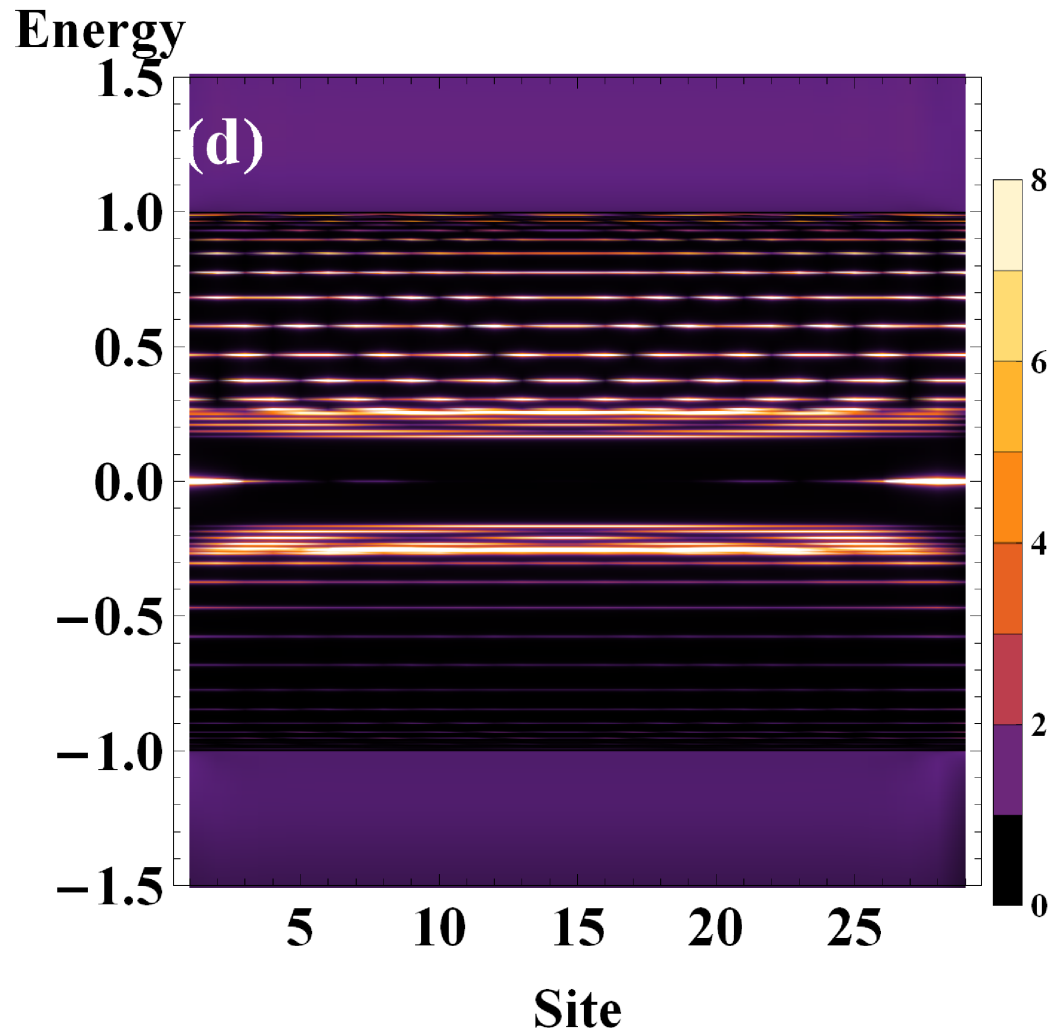
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Spectrum of a quantum wire has a series of Andreev states.

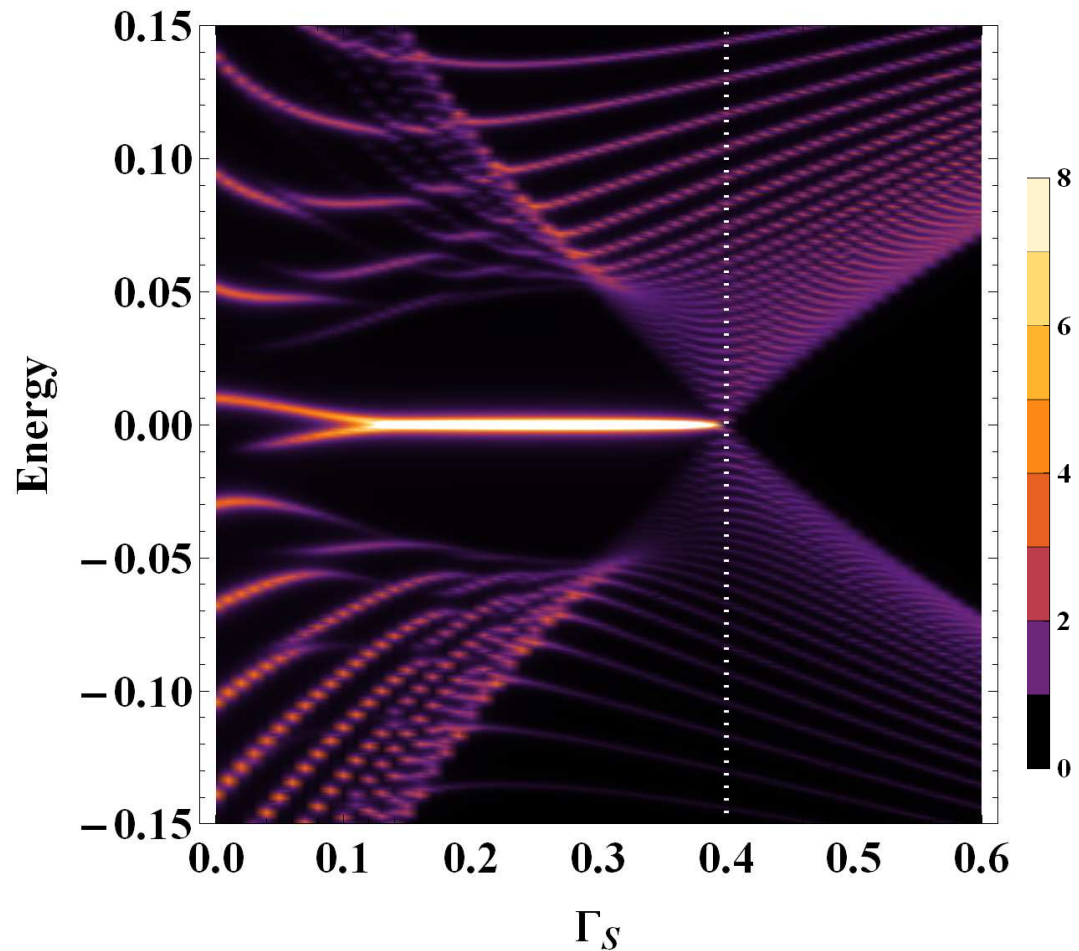
Andreev vs Majorana states

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Spin-orbit coupling can induce the Majorana-type quasiparticles.

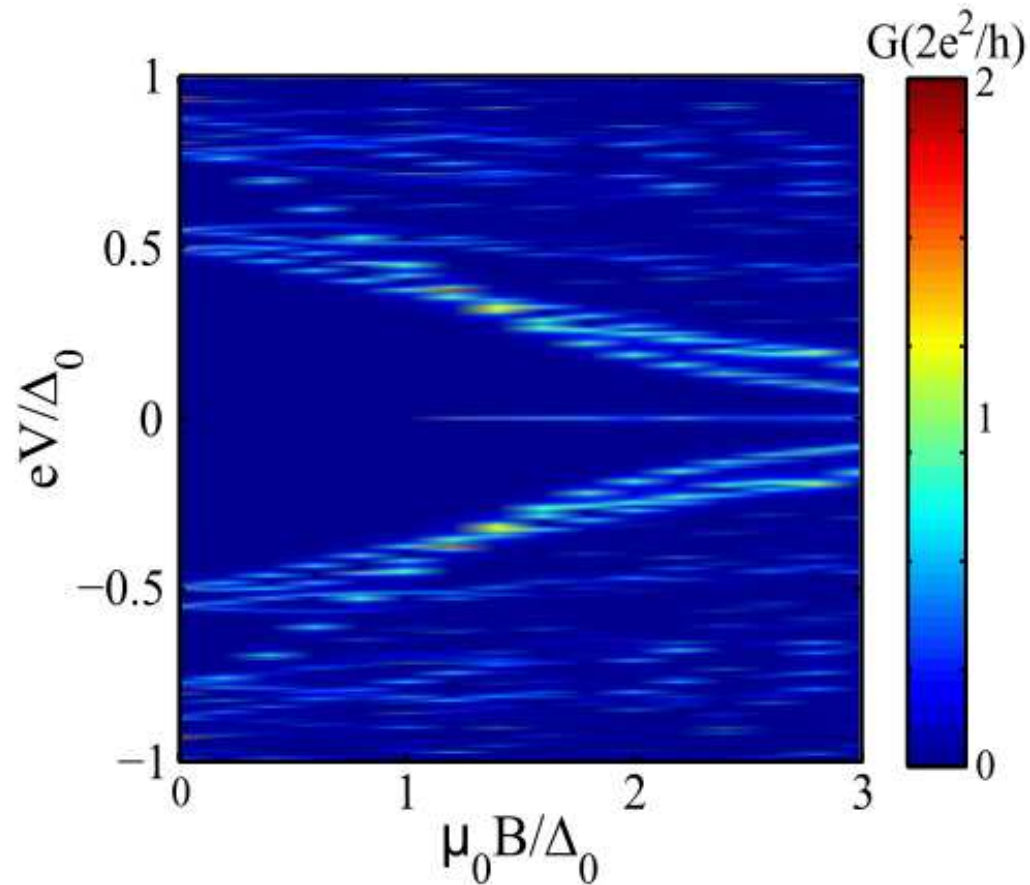
Andreev vs Majorana states – a story of mutation



Majorana quasiparticles appear at the edges of a quantum wire.

D. Chevallier, P. Simon, and C. Bena, Phys. Rev. B **88**, 165401 (2013).

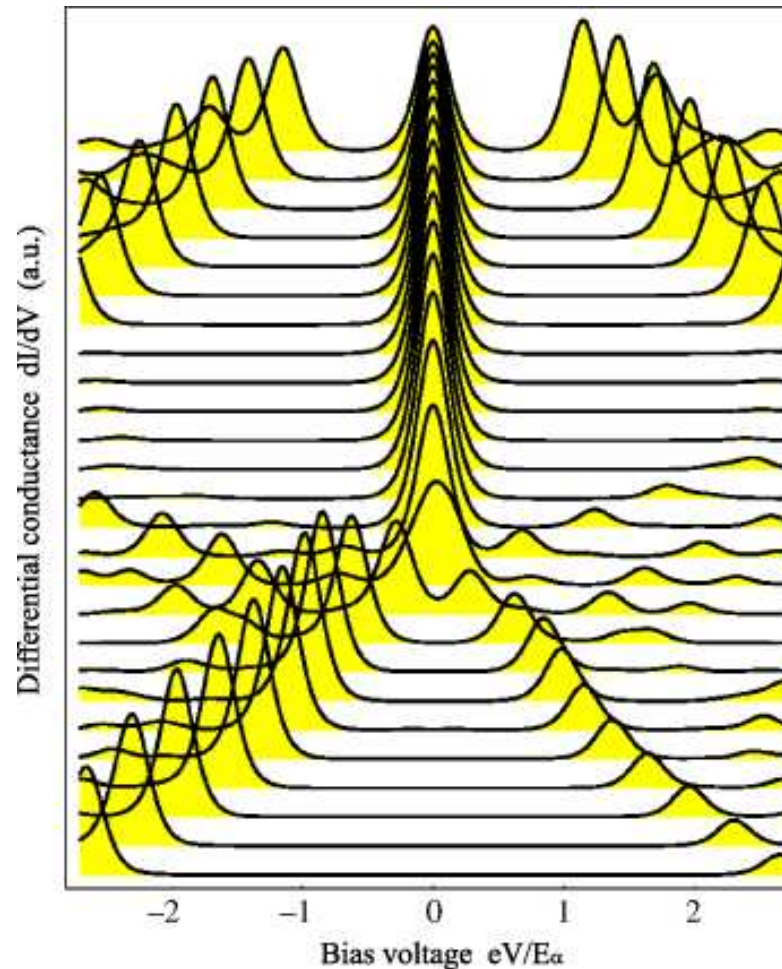
Andreev vs Majorana states – a story of mutation



Quasiparticles at the edge of a quantum wire for varying magnetic field.

J. Liu, A.C. Potter, K.T. Law, and P.A. Lee, Phys. Rev. Lett. 109, 267002 (2012).

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Quasiparticles at the edge of a quantum wire for varying magnetic field.

T.D. Stanescu, R.M. Lutchyn, and S. Das Sarma, Phys. Rev. B **84**, 144522 (2011).

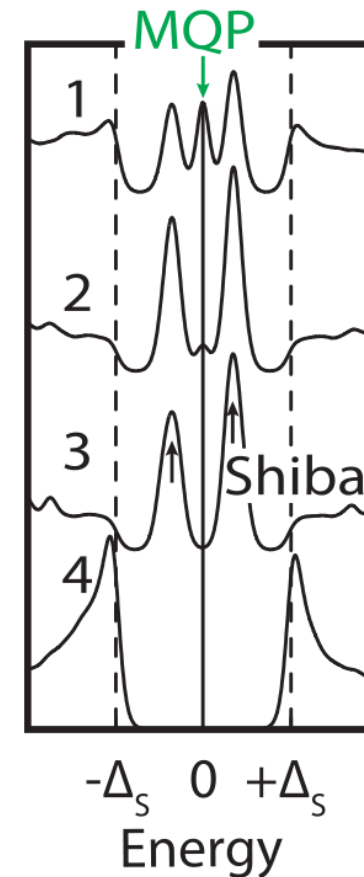
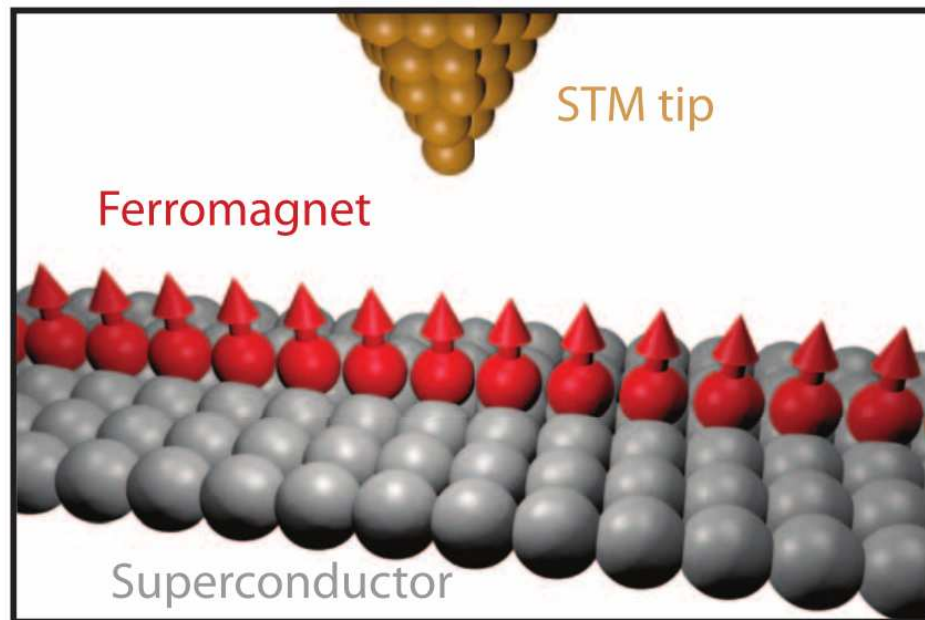
Experimental results

– **for Majorana quasiparticles**

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A chain of iron atoms deposited on a surface of superconducting lead



STM measurements provided evidence for:

⇒ **Majorana bound states at the edges of a chain.**

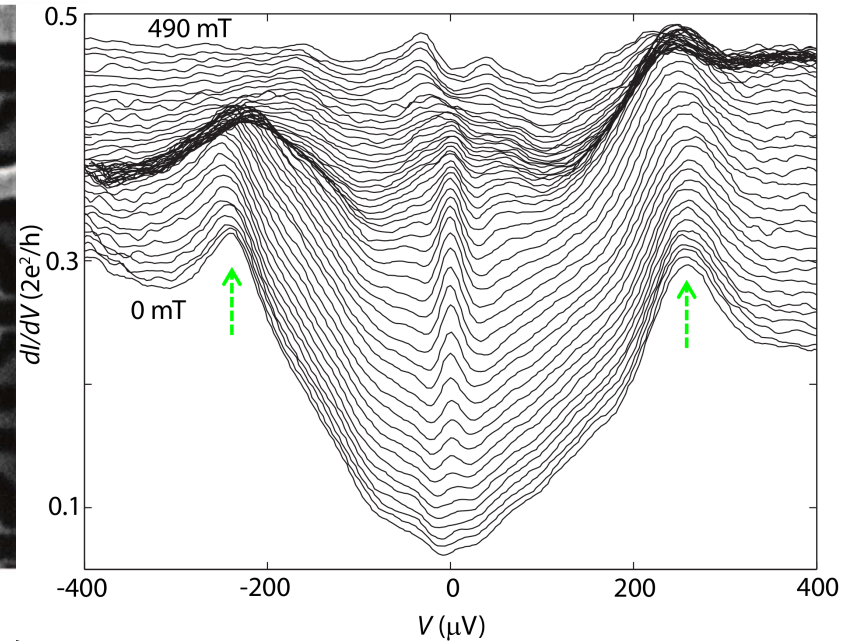
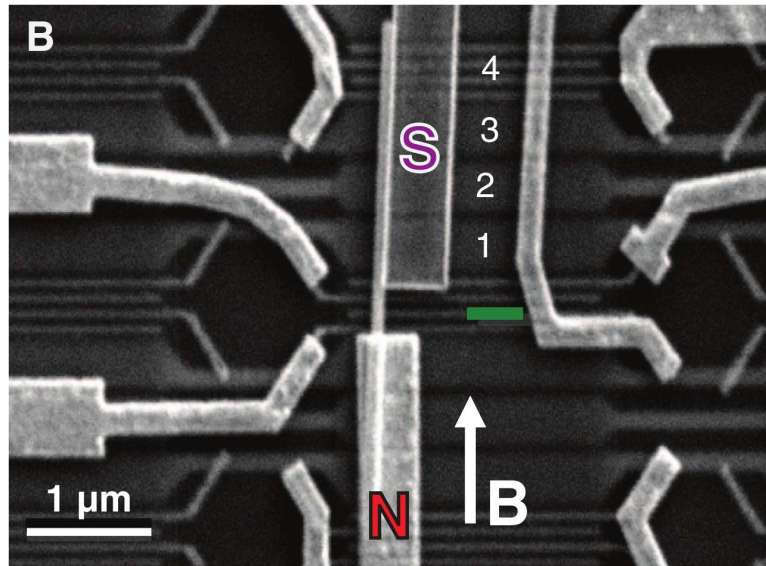
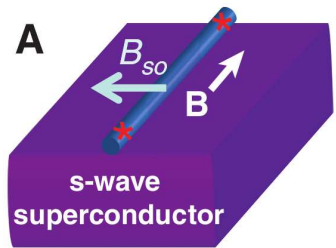
S. Nadj-Perge, ..., and A. Yazdani, Science **346**, 602 (2014).

/ Princeton University, Princeton (NJ), USA /

Experimental results

– for Majorana quasiparticles

InSb nanowire between a metal (gold) and a superconductor (Nb-Ti-N)



dI/dV measured at 70 mK for varying magnetic field B indicated:

⇒ a zero-bias enhancement due to Majorana state

V. Mourik, ..., and L.P. Kouwenhoven, Science **336**, 1003 (2012).

/ Kavli Institute of Nanoscience, Delft Univ., Netherlands /